

# **Dynamics and Time Integration**

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**Computer Graphics  
CMU 15-462/15-662**

# Last time: animation

- Added motion to our model
- Interpolate keyframes
- Still a lot of work!
- Today: physically-based animation
  - often less manual labor
  - often more compute-intensive
- Leverage tools from physics
  - dynamical descriptions
  - numerical integration
- Payoff: beautiful, complex behavior from simple models
- Widely-used techniques in modern film (and games!)



# Dynamical Description of Motion

***“A change in motion is proportional to the motive force impressed and takes place along the straight line in which that force is impressed.”***

**—Sir Isaac Newton, 1687**

***“Dynamics is concerned with the study of forces and their effect on motion, as opposed to kinematics, which studies the motion of objects without reference to its causes.”***

**—Sir Wiki Pedia, 2015**

**(Q: Is keyframe interpolation *dynamic*, or *kinematic*?)**

# The Animation Equation

- Already saw the *rendering equation*
- What's the *animation equation*?

$$F = ma$$

The diagram shows the equation  $F = ma$  in a serif font. Three red arrows point from labels to the variables: 'force' points to  $F$ , 'mass' points to  $m$ , and 'acceleration' points to  $a$ .



# The “Animation Equation,” revisited

■ Well actually there are some more equations...

■ Let's be more careful:

- Any system has a *configuration*  $q(t)$

- It also has a *velocity*  $\dot{q} := \frac{d}{dt} q$

- And some kind of *mass*  $M$

- There are probably some *forces*  $F$

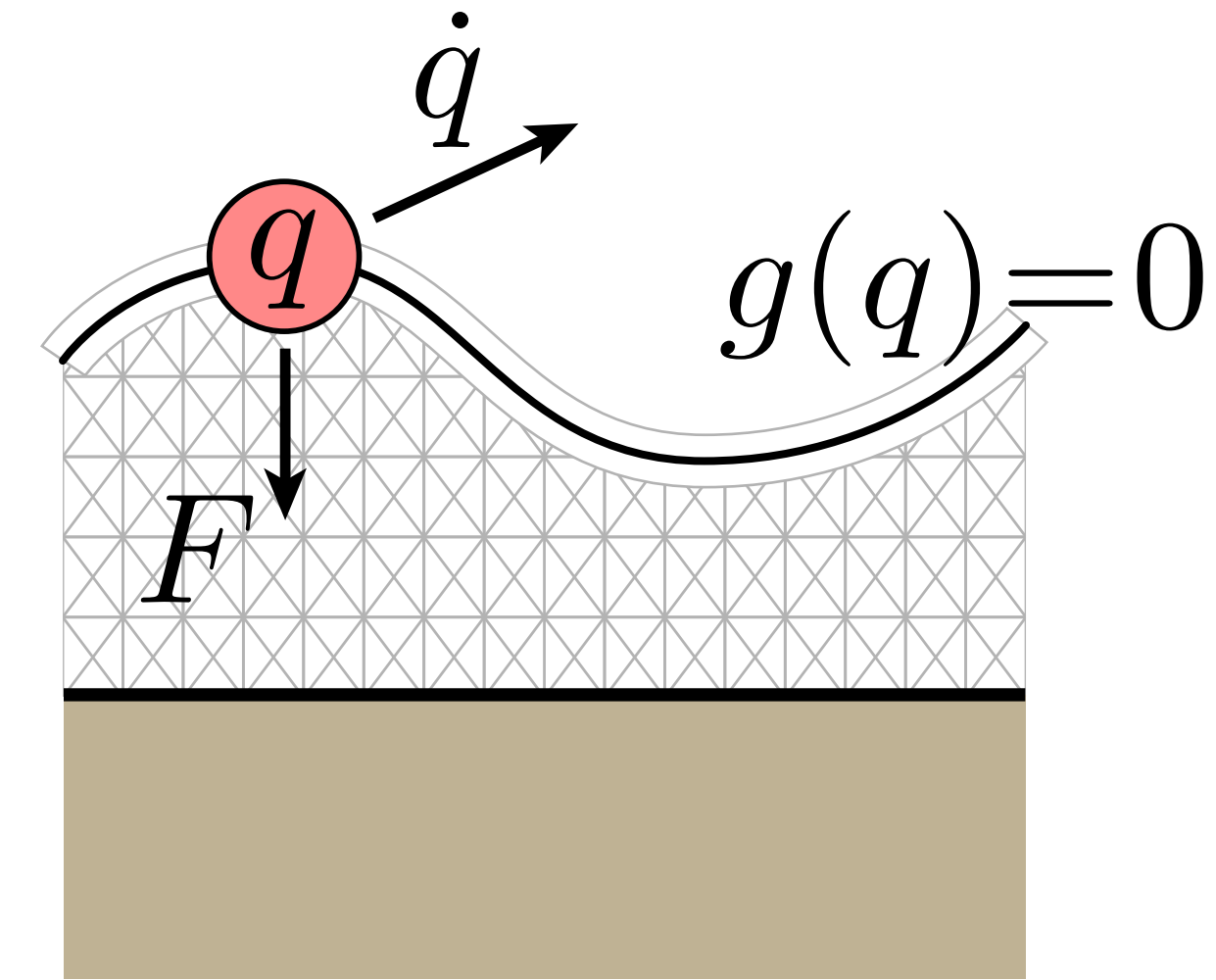
- And also some *constraints*  $g(q, \dot{q}, t) = 0$

■ E.g., could write Newton's 2nd law as  $\ddot{q} = F/m$

■ Makes two things clear:

- acceleration is 2nd time derivative of configuration

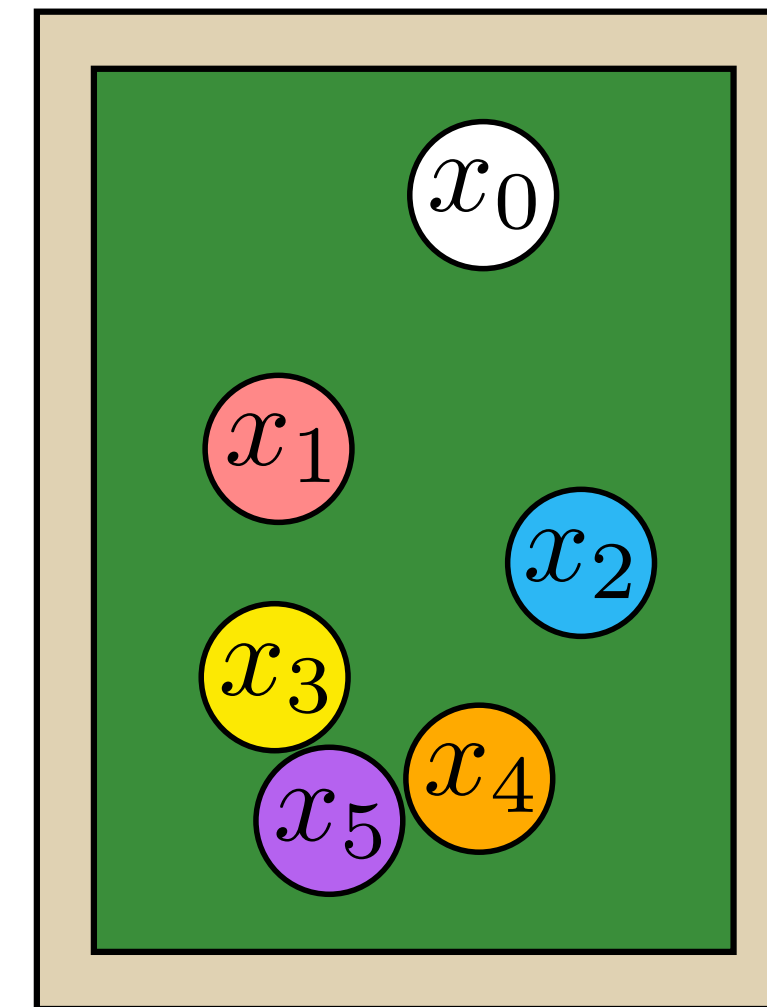
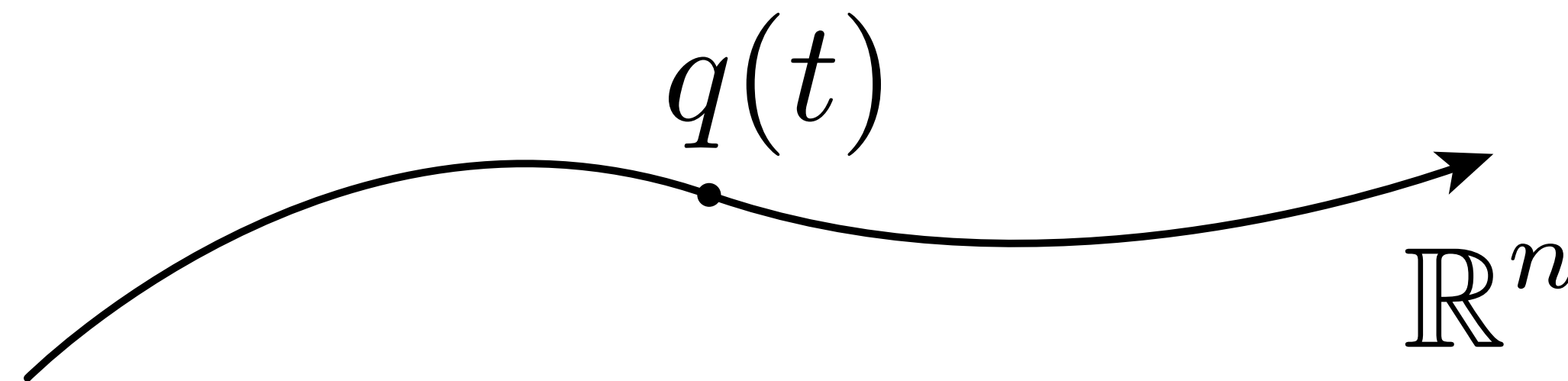
- ultimately, we want to solve for the configuration  $q$



# Generalized Coordinates

- Often describing systems with many, many moving pieces
- E.g., a collection of billiard balls, each with position  $x_i$
- Collect them all into a single vector of *generalized coordinates*:

$$q = (x_0, x_1, \dots, x_n)$$

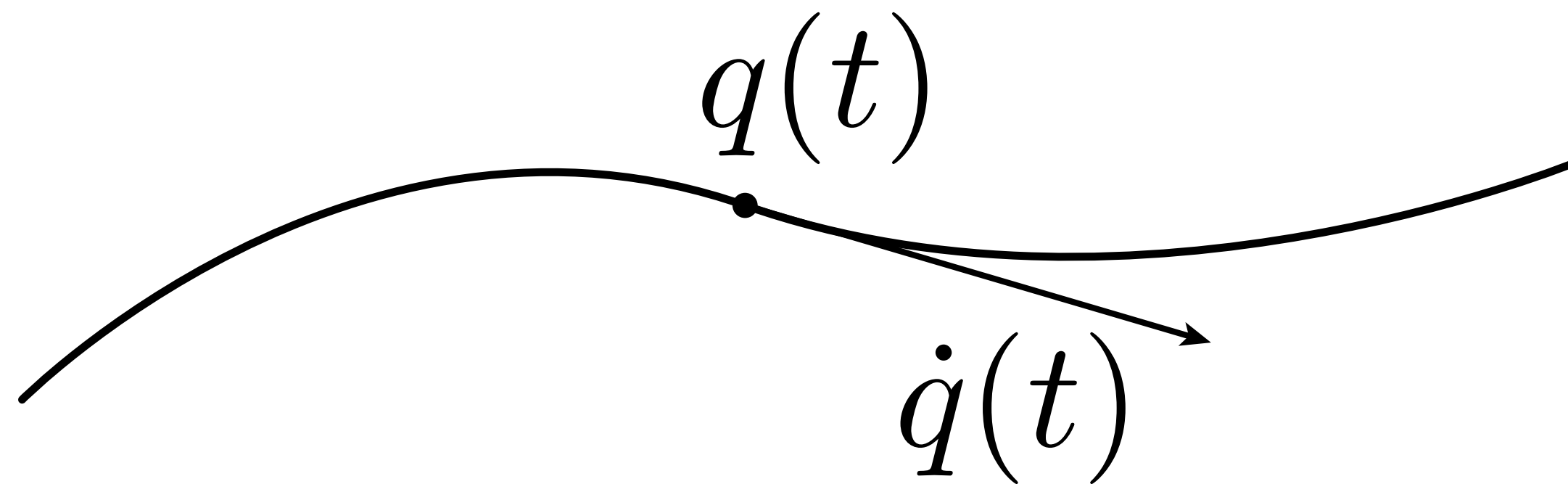


- Can think of  $q$  as a *single point* moving along a trajectory in  $R^n$
- This way of thinking naturally maps to the way we actually solve equations on a computer: all variables are often “stacked” into a big long vector and handed to a solver.
- (...So why not write things down this way in the first place?)

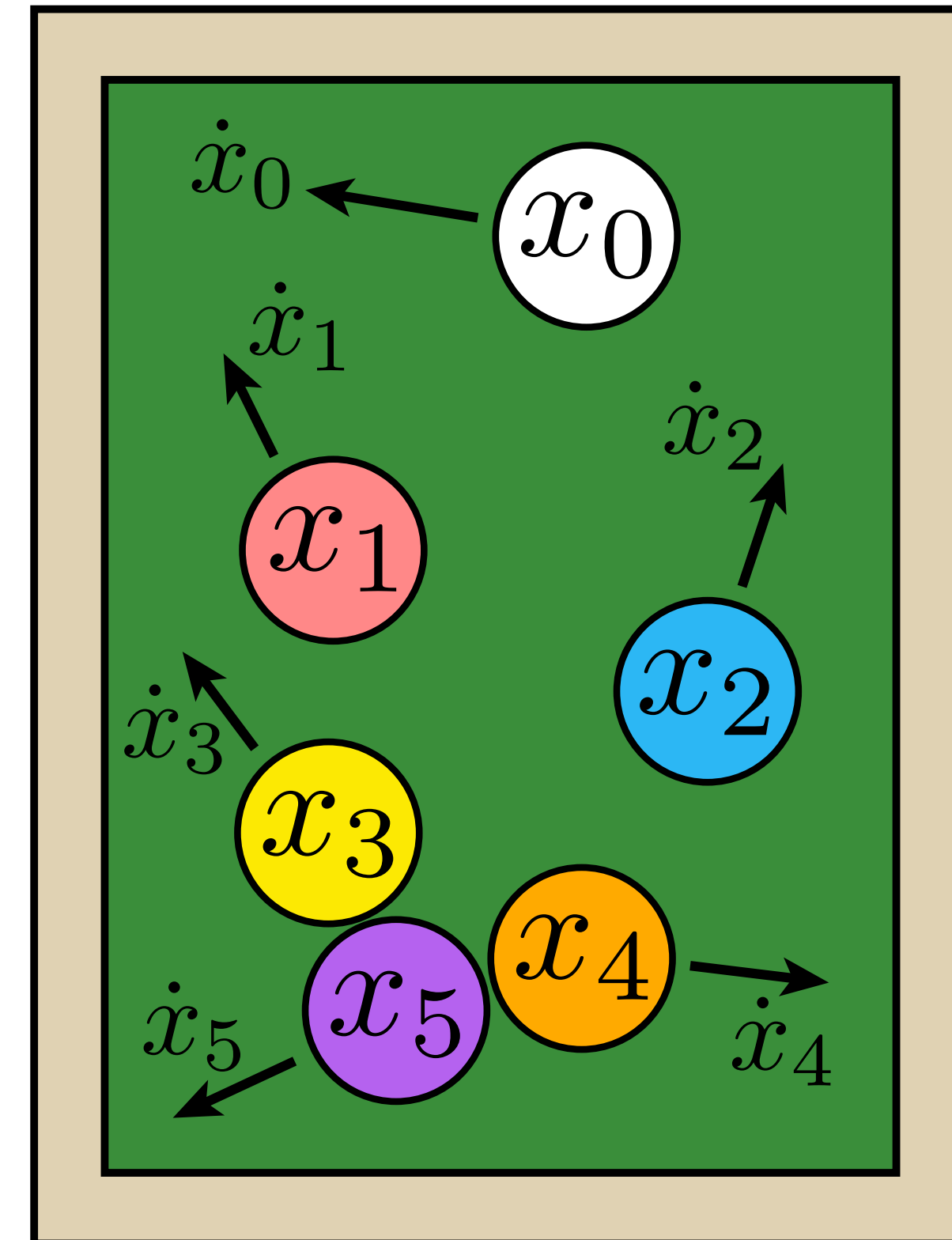
# Generalized Velocity

- Not much more to say about generalized velocity: it's the time derivative of the generalized coordinates!

$$\dot{q} = (\dot{x}_0, \dot{x}_1, \dots, \dot{x}_n)$$



**All of life (and physics) is just traveling along a curve...**



# Ordinary Differential Equations

- Many dynamical systems can be described via an *ordinary differential equation (ODE)* in generalized coordinates:

change in configuration over time velocity function

$$\frac{d}{dt} q = f(q, \dot{q}, t)$$

- ODE doesn't have to describe mechanical phenomenon, e.g.,

$$\frac{d}{dt} u(t) = au$$

“rate of growth is proportional to value”

- **Solution?**  $u(t) = be^{at}$
- Describes exponential decay ( $a < 1$ ), or really great stock ( $a > 1$ )
- “Ordinary” means “involves derivatives in time but not space”
- We'll talk about spatial derivatives (PDEs) in another lecture...

# Dynamics via ODEs

- Another key example: Newton's 2nd law!

$$\ddot{q} = F/m$$

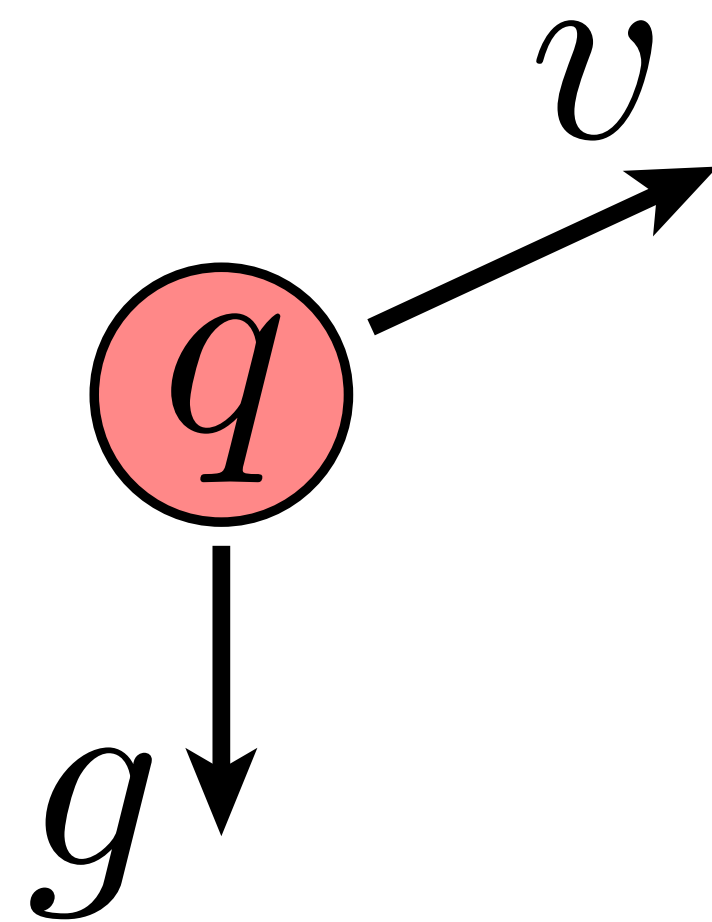
- "Second order" ODE since we take *two* time derivatives
- Can also write as a *system* of two *first order* ODEs, by introducing new "dummy" variable for velocity:

$$\begin{aligned}\dot{q} &= v \\ \dot{v} &= F/m\end{aligned}$$

- Splitting things up this way will make it easy to talk about solving these equations numerically (among other things)

# Simple Example: Throwing a Rock

- Consider a rock\* of mass  $m$  tossed under force of gravity  $g$
- Easy to write dynamical equations, since only force is gravity:



$$\ddot{q} = g/m \quad \text{or} \quad \begin{aligned} \dot{q} &= v \\ \dot{v} &= g/m \end{aligned}$$

**Solution:**

$$\begin{aligned} v(t) &= v_0 + \frac{t}{m}g \\ q(t) &= q_0 + tv_0 + \frac{t^2}{2m}g \end{aligned}$$



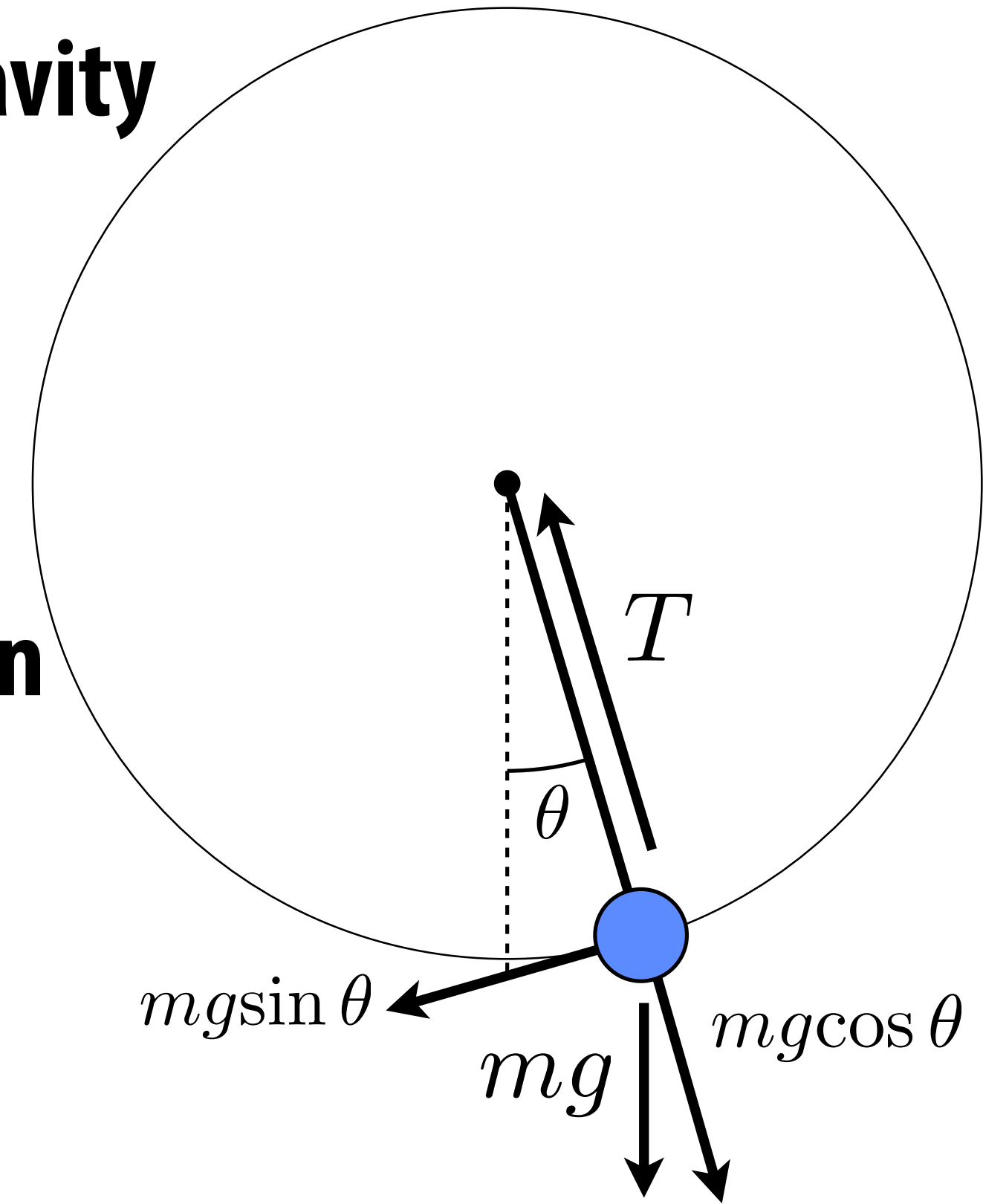
**(What do we need a computer for?!)**

\*Yes, this rock is spherical and has uniform density.

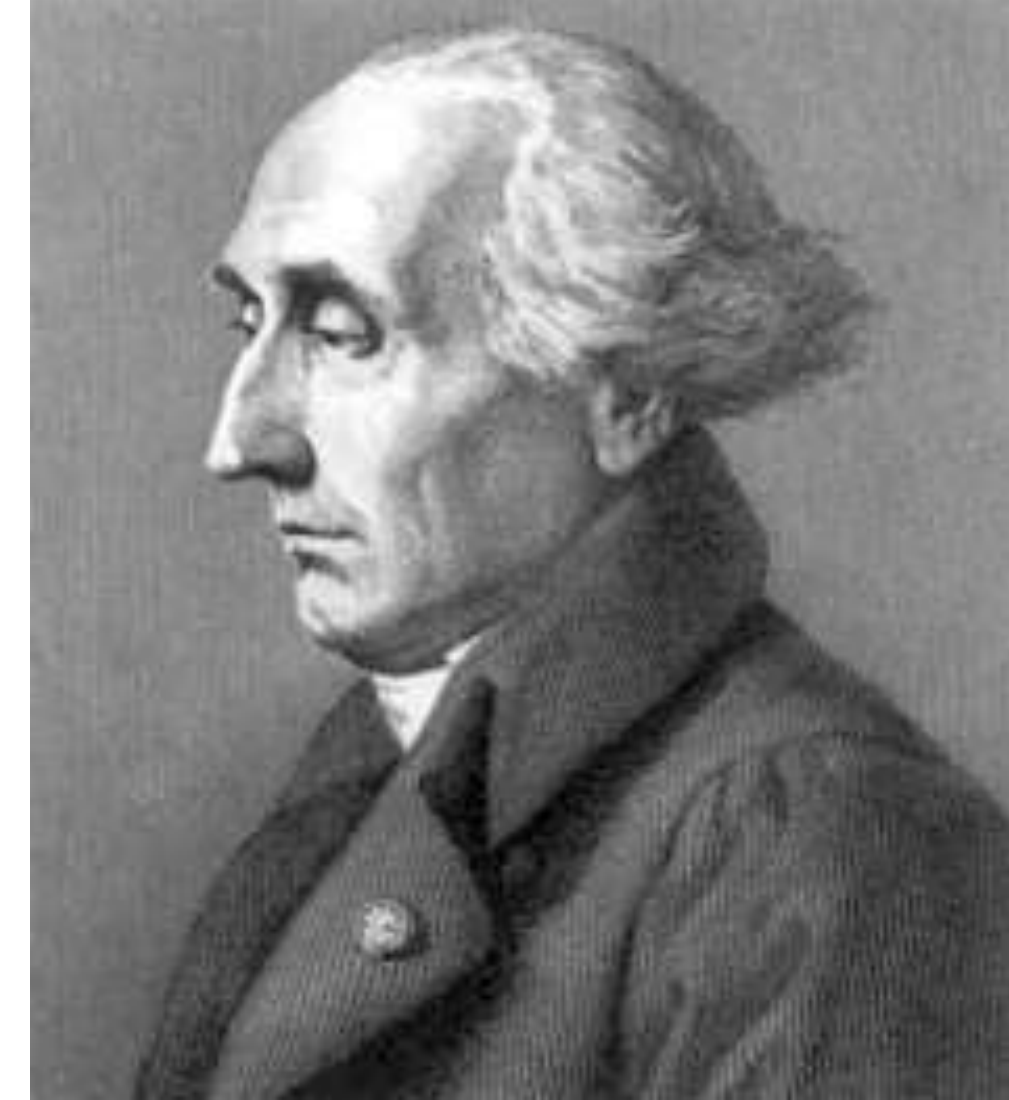


# Slightly Harder Example: Pendulum

- Mass on end of a bar, swinging under gravity
- What are the equations of motion?
- Same as “rock” problem, but *constrained*
- Could use a “*force diagram*”
  - You probably did this for many hours in high school/college
  - Let's do something new & different!



# Lagrangian Mechanics



Joe Lagrange

## ■ Beautifully simple recipe:

1. Write down kinetic energy  $K$
2. Write down potential energy  $U$
3. Write down *Lagrangian*  $\mathcal{L} := K - U$
4. Dynamics then given by *Euler-Lagrange equation*

$$\begin{array}{c} \text{becomes (generalized)} \\ \text{"MASS TIMES ACCELERATION"} \end{array} \longrightarrow \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} = \frac{\partial \mathcal{L}}{\partial q} \longleftarrow \begin{array}{c} \text{becomes (generalized) "FORCE"} \end{array}$$

## ■ Why is this useful?

- often easier to come up with (scalar) energies than forces
- very general, works in any kind of generalized coordinates
- helps develop nice class of numerical integrators (symplectic)

Great reference: Sussman & Wisdom, "Structure and Interpretation of Classical Mechanics"

# Lagrangian Mechanics - Example

- Generalized coordinates for pendulum?

$$q = \theta \leftarrow \text{just one coordinate: angle with the vertical direction}$$

- Kinetic energy (mass  $m$ )?

$$K = \frac{1}{2} I \omega^2 = \frac{1}{2} m L^2 \dot{\theta}^2$$

- Potential energy?

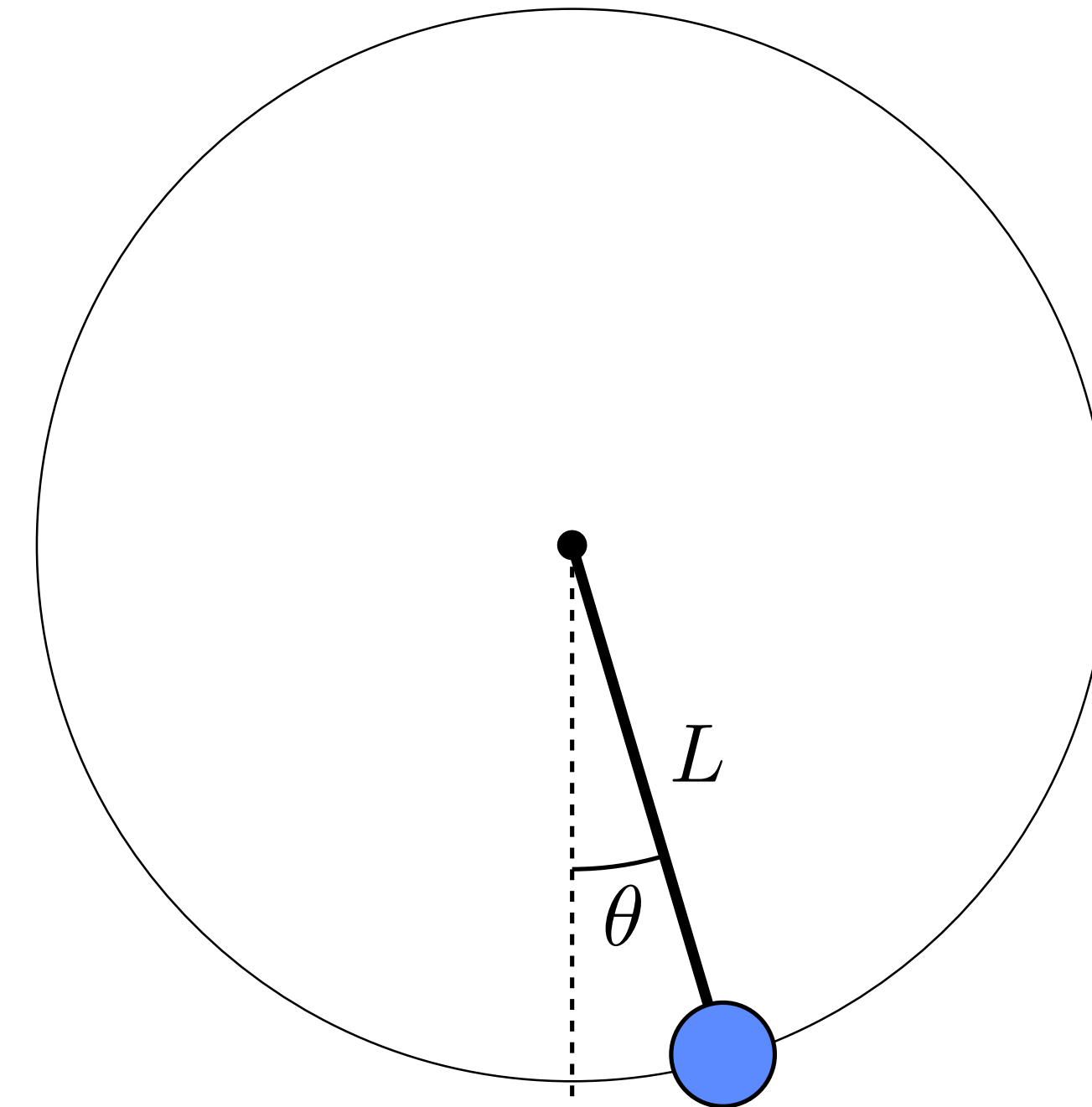
$$U = mgh = -mgL \cos \theta$$

- Euler-Lagrange equations? (from here, just “plug and chug”—even a computer could do it!)

$$\mathcal{L} = K - U = m \left( \frac{1}{2} L^2 \dot{\theta}^2 + gL \cos \theta \right)$$

$$\frac{\partial \mathcal{L}}{\partial \dot{q}} = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = mL^2 \dot{\theta} \qquad \frac{\partial \mathcal{L}}{\partial q} = \frac{\partial \mathcal{L}}{\partial \theta} = -mgL \sin \theta$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} = \frac{\partial \mathcal{L}}{\partial q} \quad \Rightarrow \quad \boxed{\ddot{\theta} = -\frac{g}{L} \sin \theta}$$



# Solving the Pendulum

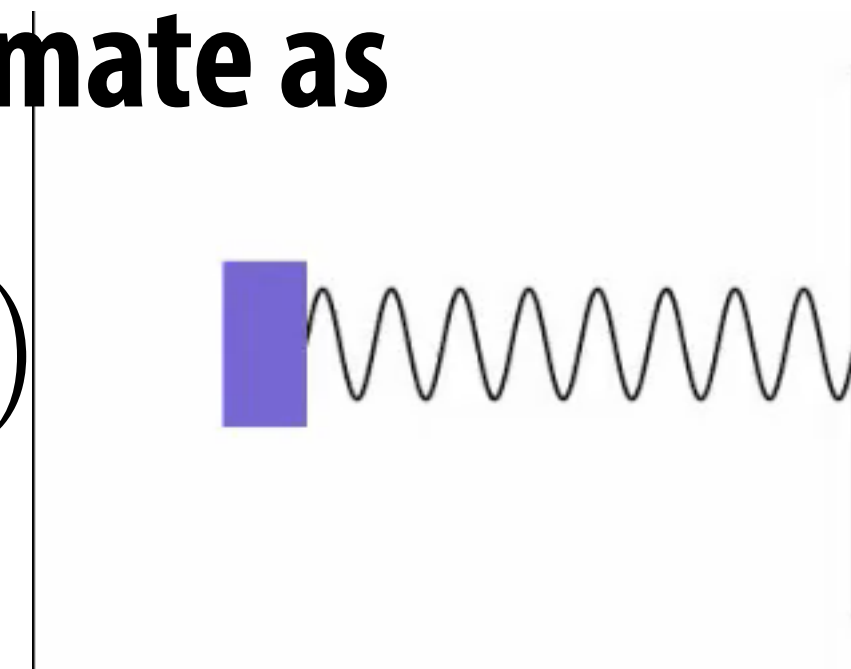
- Great, now we have a nice simple equation for the pendulum:

$$\ddot{\theta} = -\frac{g}{L} \sin \theta$$

- For small angles (e.g., clock pendulum) can approximate as

$$\ddot{\theta} = -\frac{g}{L} \theta \quad \Rightarrow \quad \theta(t) = a \cos(t \sqrt{g/L} + b)$$

**“harmonic oscillator”**

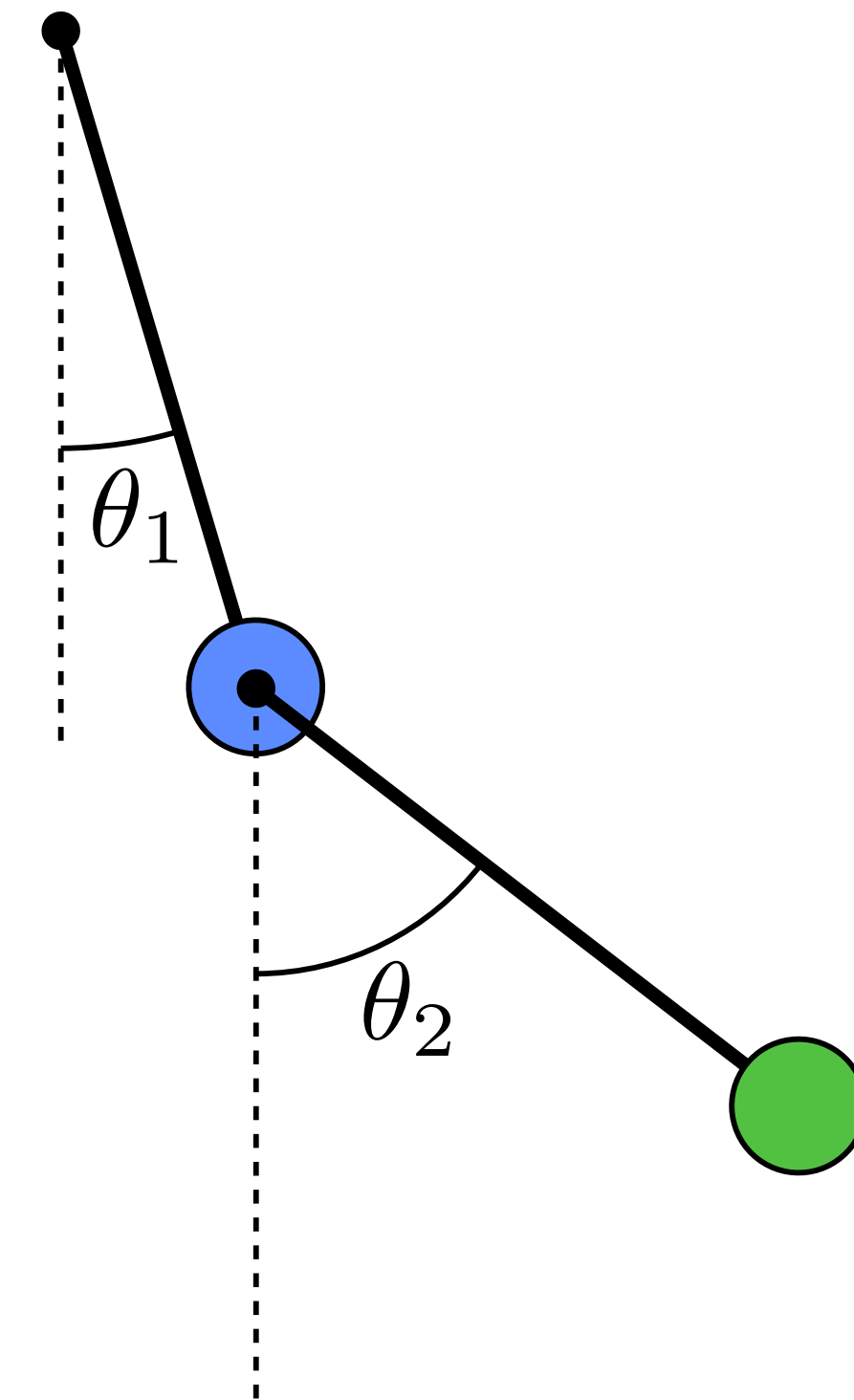
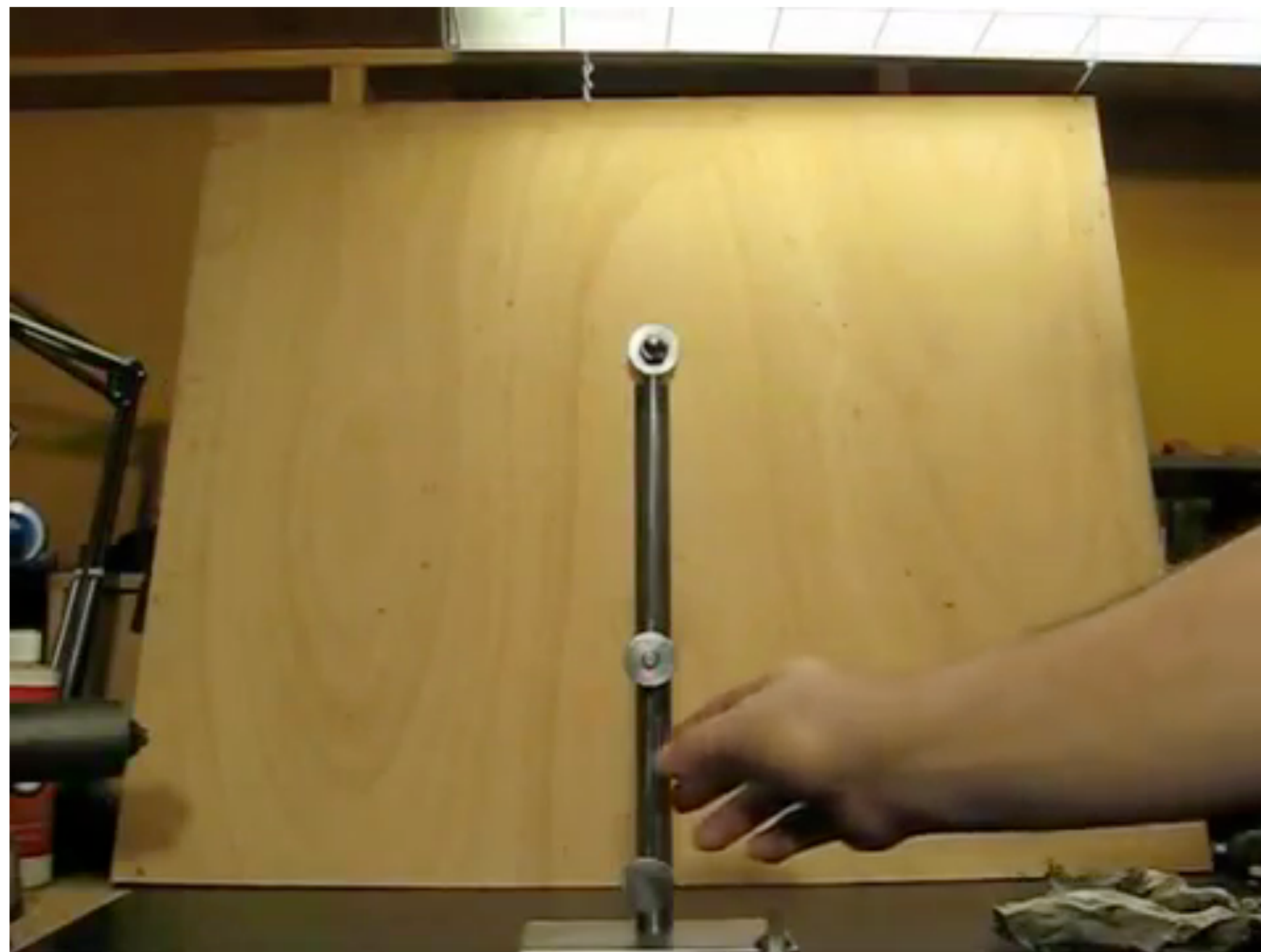


- In general, there is *no closed form solution!*
- Hence, we *must* use a numerical approximation
- ...And this was (almost) the simplest system we can think of!
- (What if we want to animate something more interesting?)



# Not-So-Simple Example: Double Pendulum

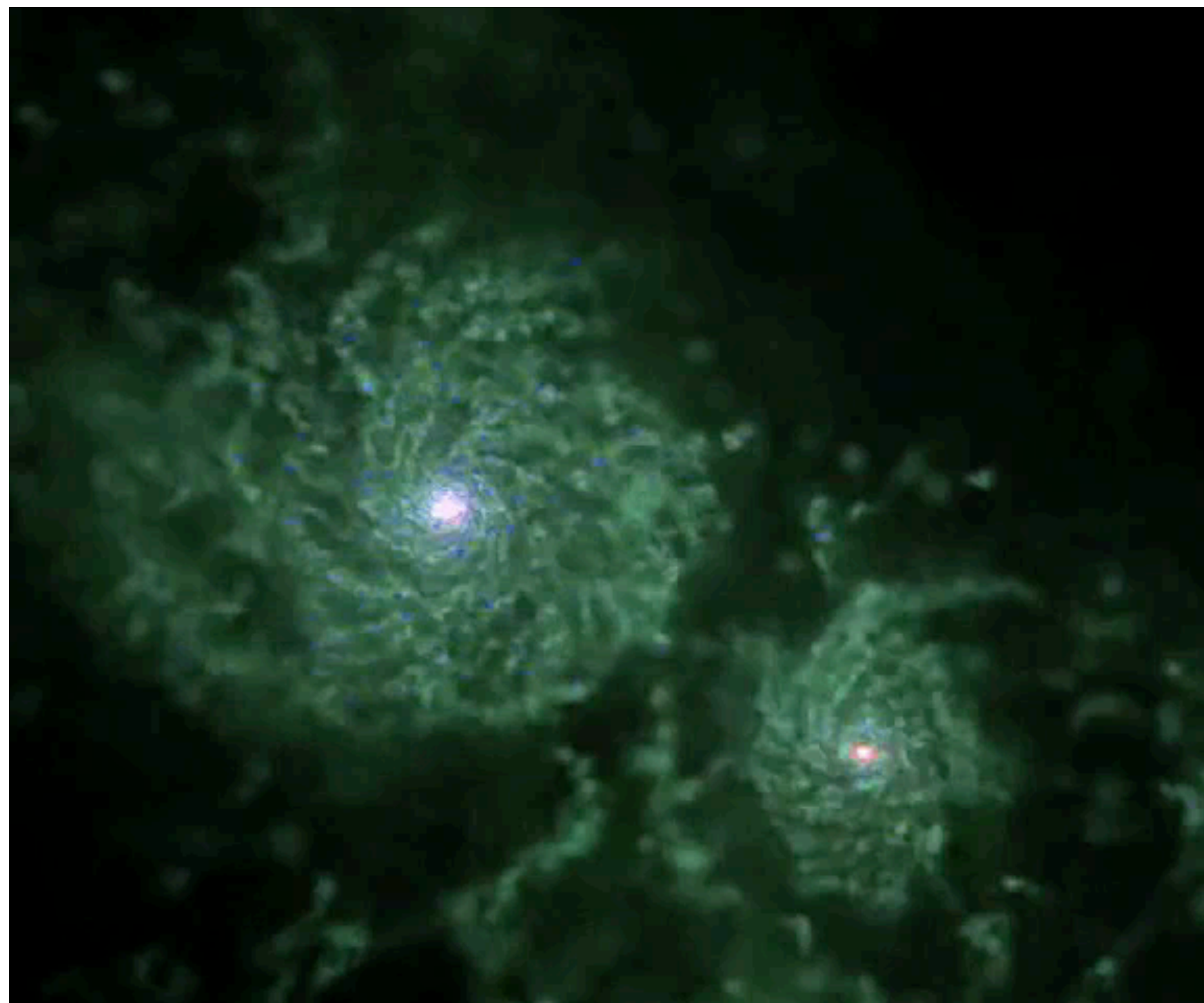
- Blue ball swings from fixed point; green ball swings from blue one
- Simple system... not-so-simple motion!
- Chaotic: perturb input, wild changes to output
- Must again use numerical approximation



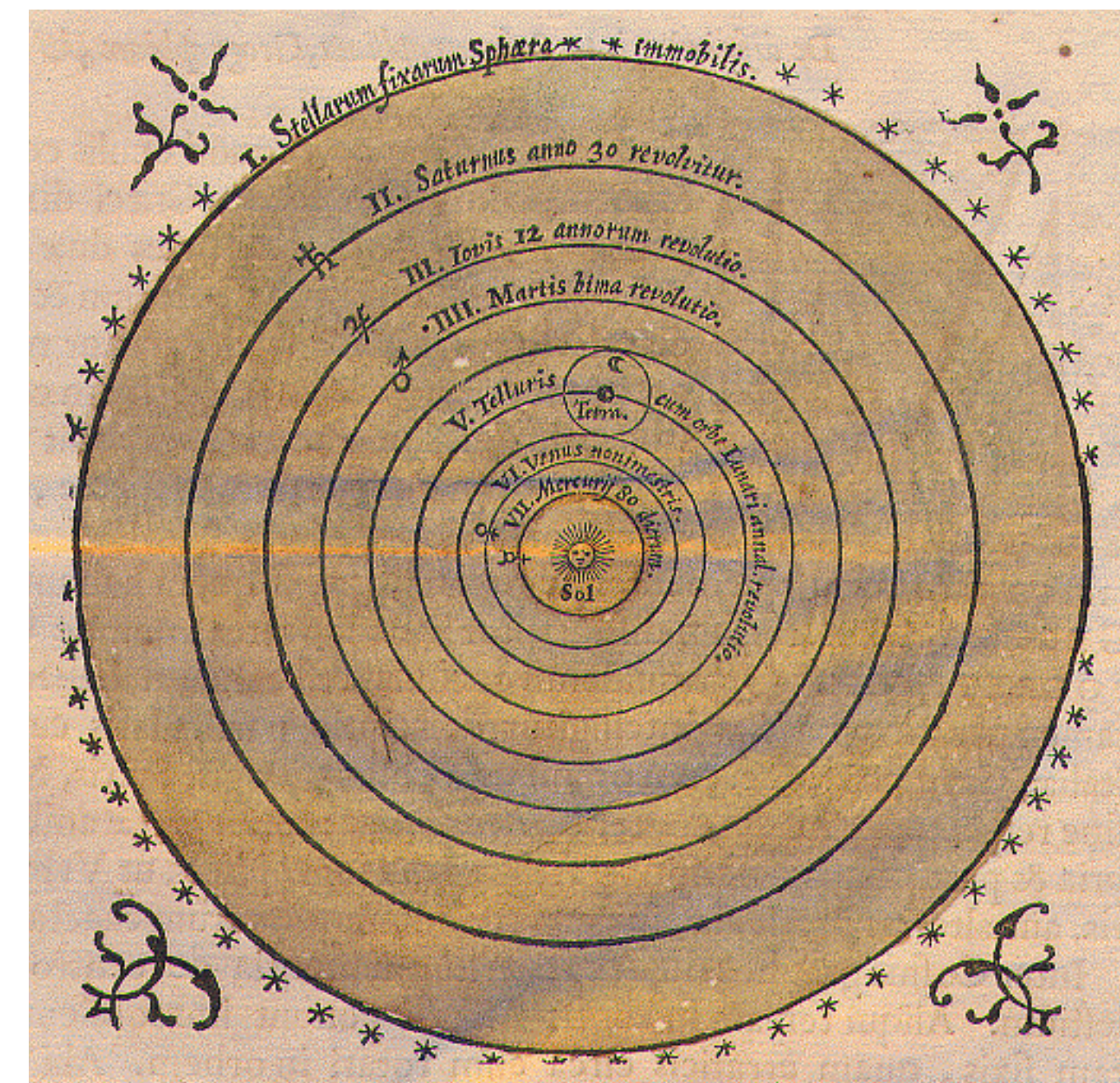


# Not-So-Simple Example: $n$ -Body Problem

- Consider the Earth, moon, and sun—where do they go?
- Solution is trivial for two bodies (e.g., assume one is fixed)
- As soon as  $n \geq 3$ , again get chaotic solutions (no closed form)
- What if we want to simulate entire *galaxies*?



Credit: Governato et al / NASA





**For animation, we *want* to simulate  
these kinds of phenomena!**

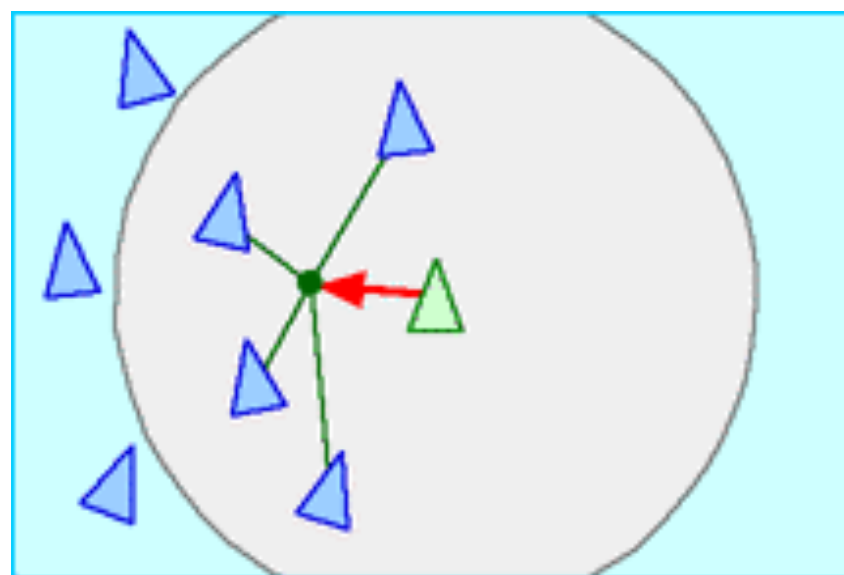
# Example: Flocking



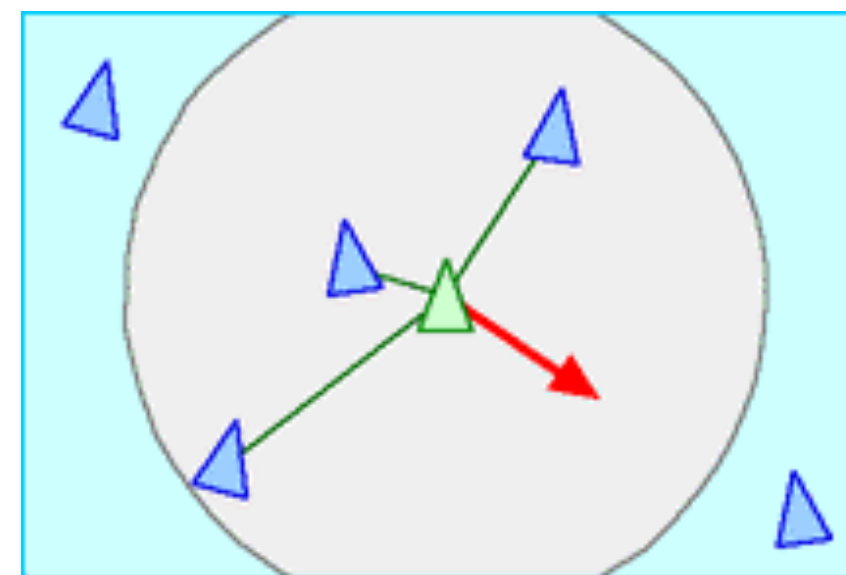


# Simulated Flocking as an ODE

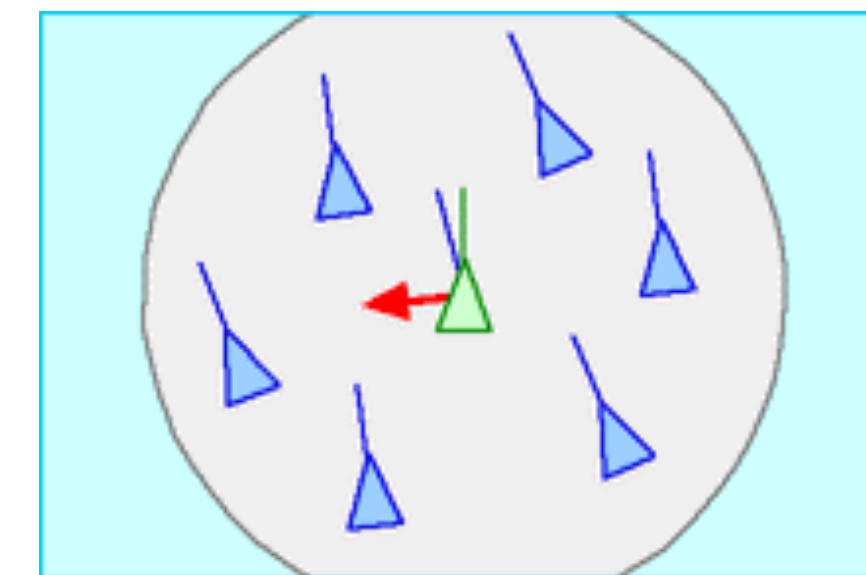
- Each bird is a particle
- Subject to very simple forces:
  - *attraction* to center of neighbors
  - *repulsion* from individual neighbors
  - *alignment* toward average trajectory of neighbors
- Solve large system of ODEs (numerically!)
- Emergent complex behavior (also seen in fish, bees, ...)



attraction



repulsion

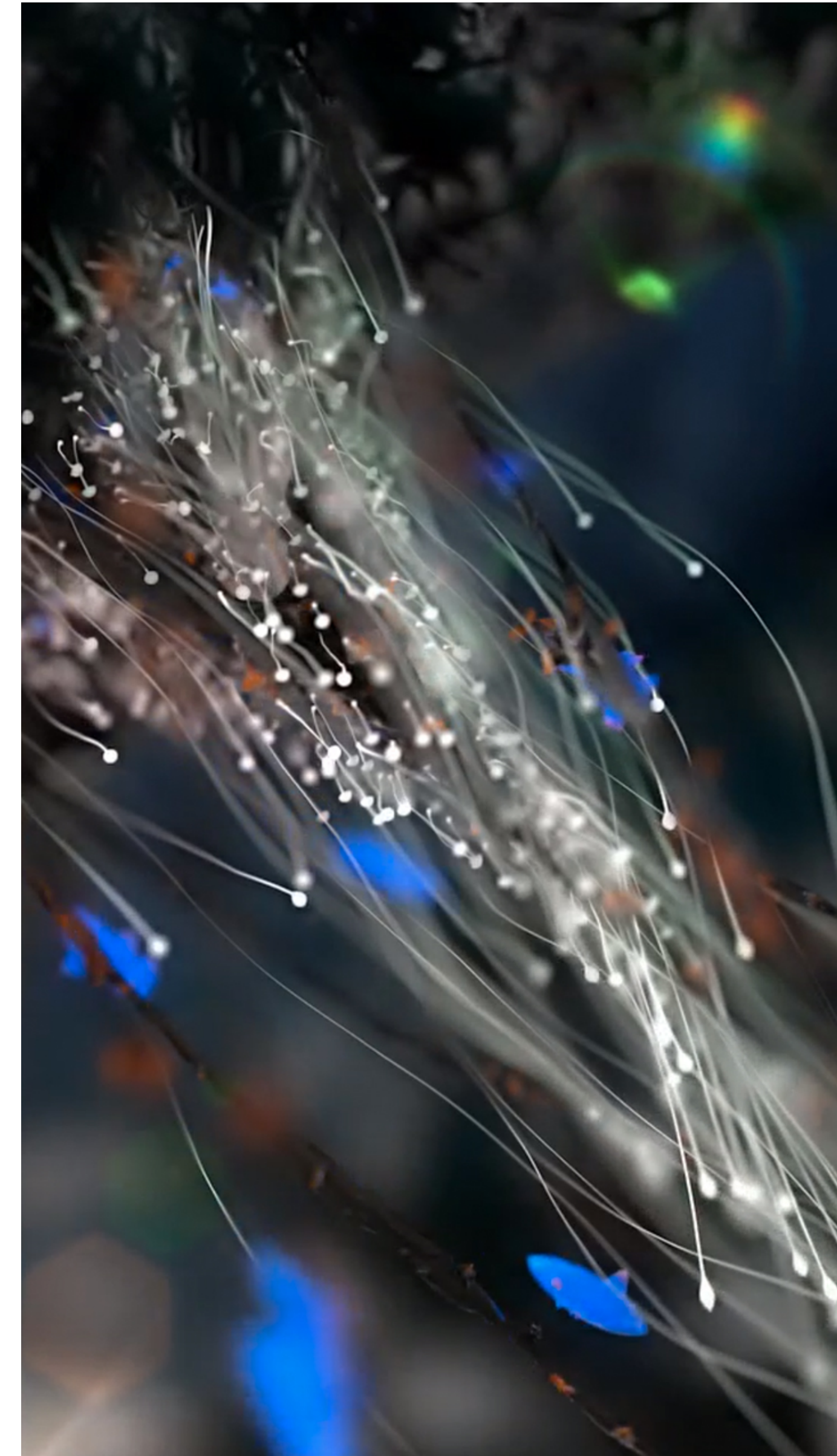


alignment



# Particle Systems

- **More generally, model phenomena as large collection of particles**
- **Each particle has a behavior described by (physical or *non-physical*) forces**
- **Extremely common in graphics/games**
  - **easy to understand**
  - **simple equation for each particle**
  - **easy to scale up/down**
- **May need *many* particles to capture certain phenomena (e.g., fluids)**
  - **may require fast hierarchical data structure (kd-tree, BVH, ...)**
  - **often better to use continuum model**





# Example: Crowds



**Where are the bottlenecks in a building plan?**



# Example: Crowds + “Rock” Dynamics



Dave Fothergill vfx



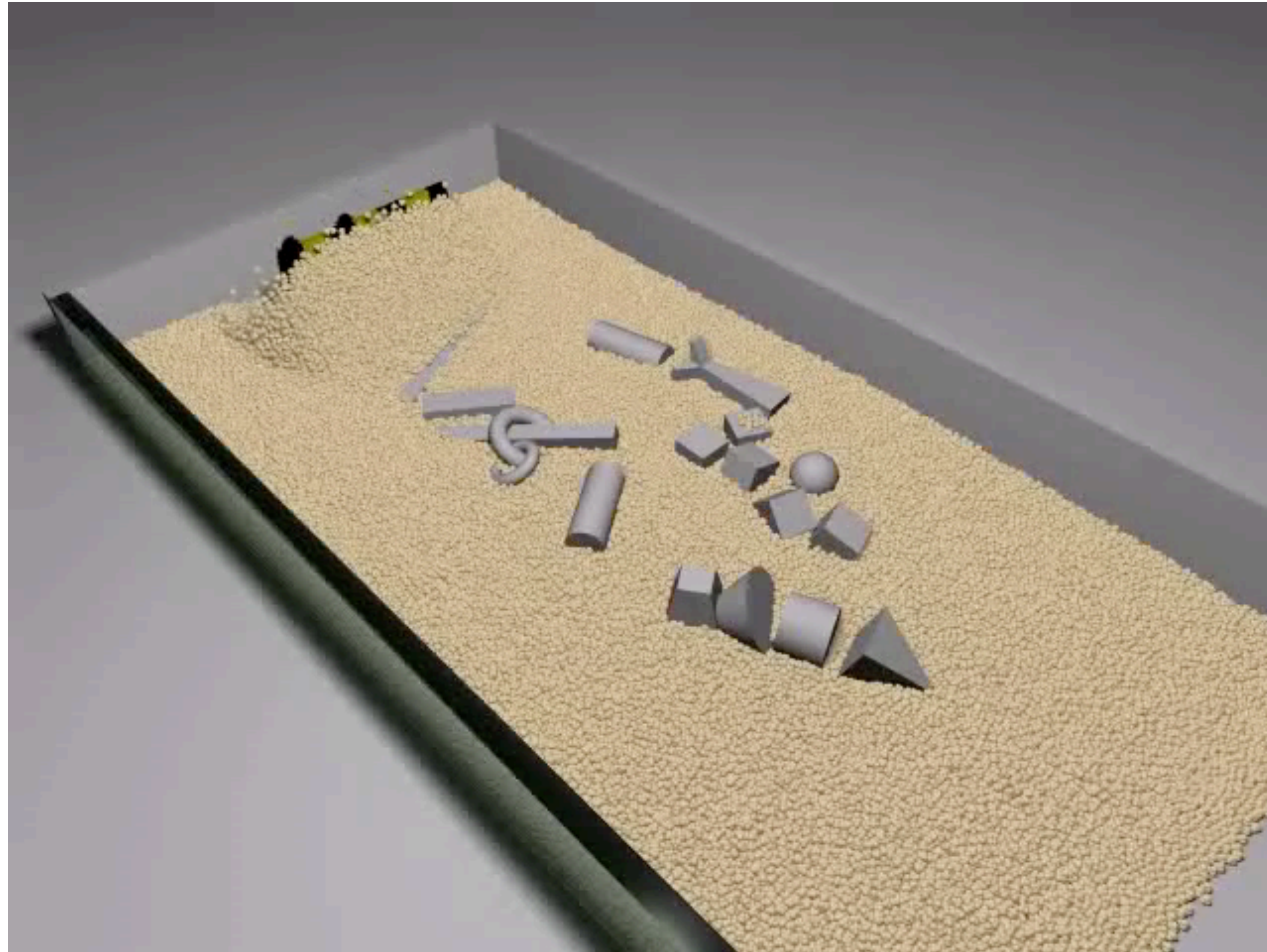
# Example: Particle-Based Fluids



**(Fluid: particles or continuum?)**



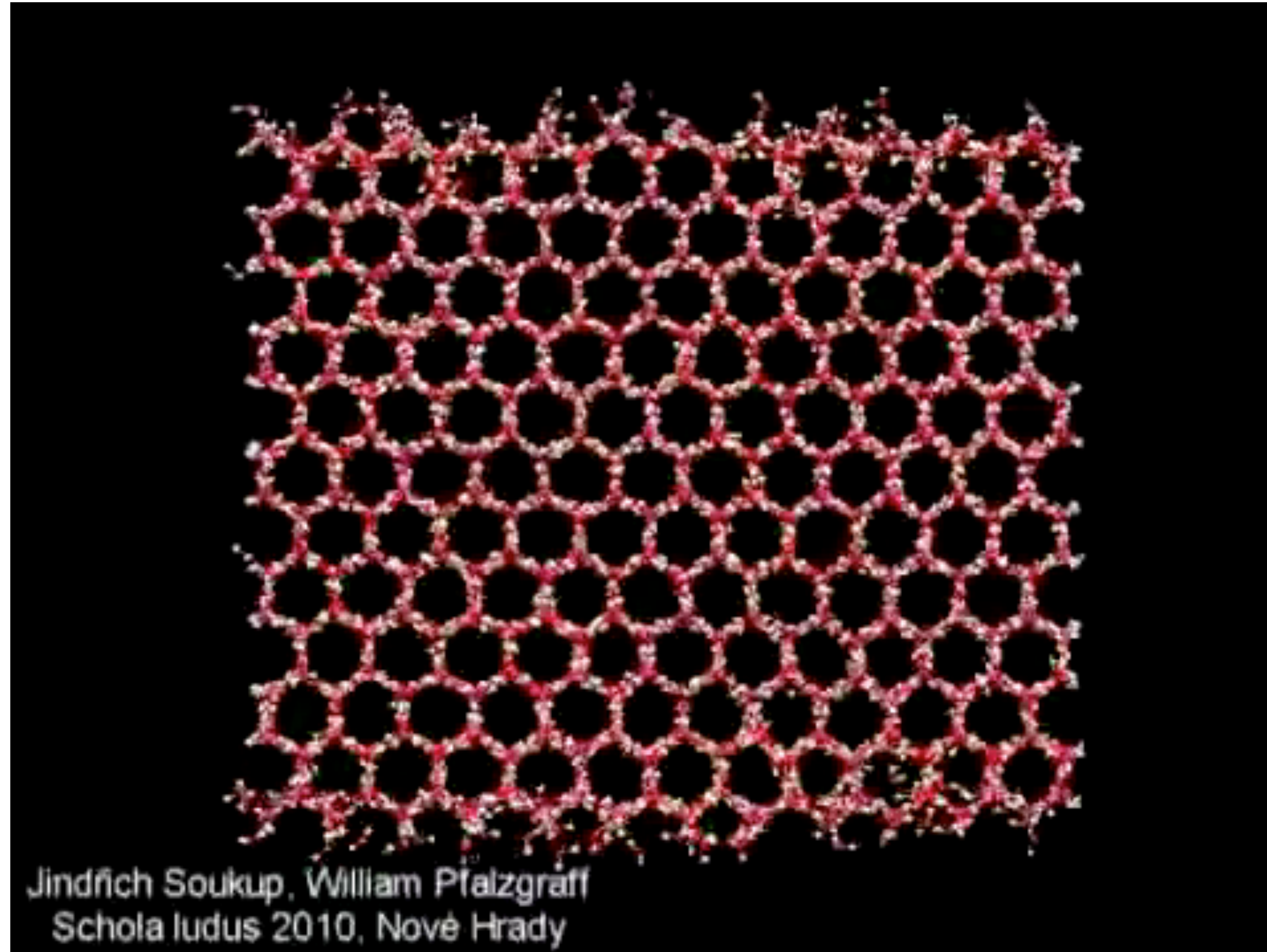
# Example: Granular Materials



**Bell et al, "Particle-Based Simulation of Granular Materials"**



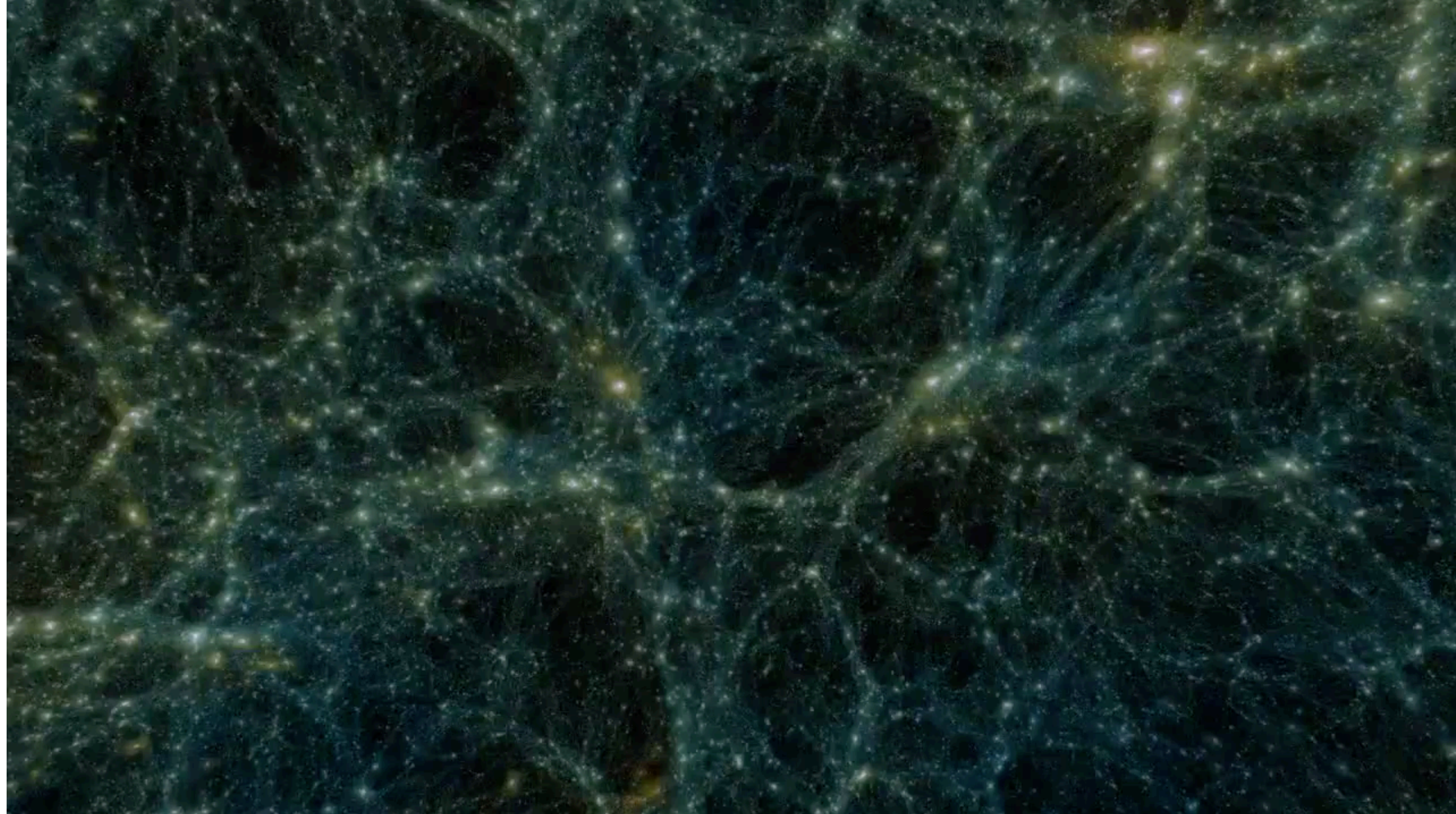
# Example: Molecular Dynamics



**(model of melting ice crystal)**



# Example: Cosmological Simulation



**Tomoaki et al -  $v^2$ GC simulation of dark matter ( $\sim 1$  trillion particles)**



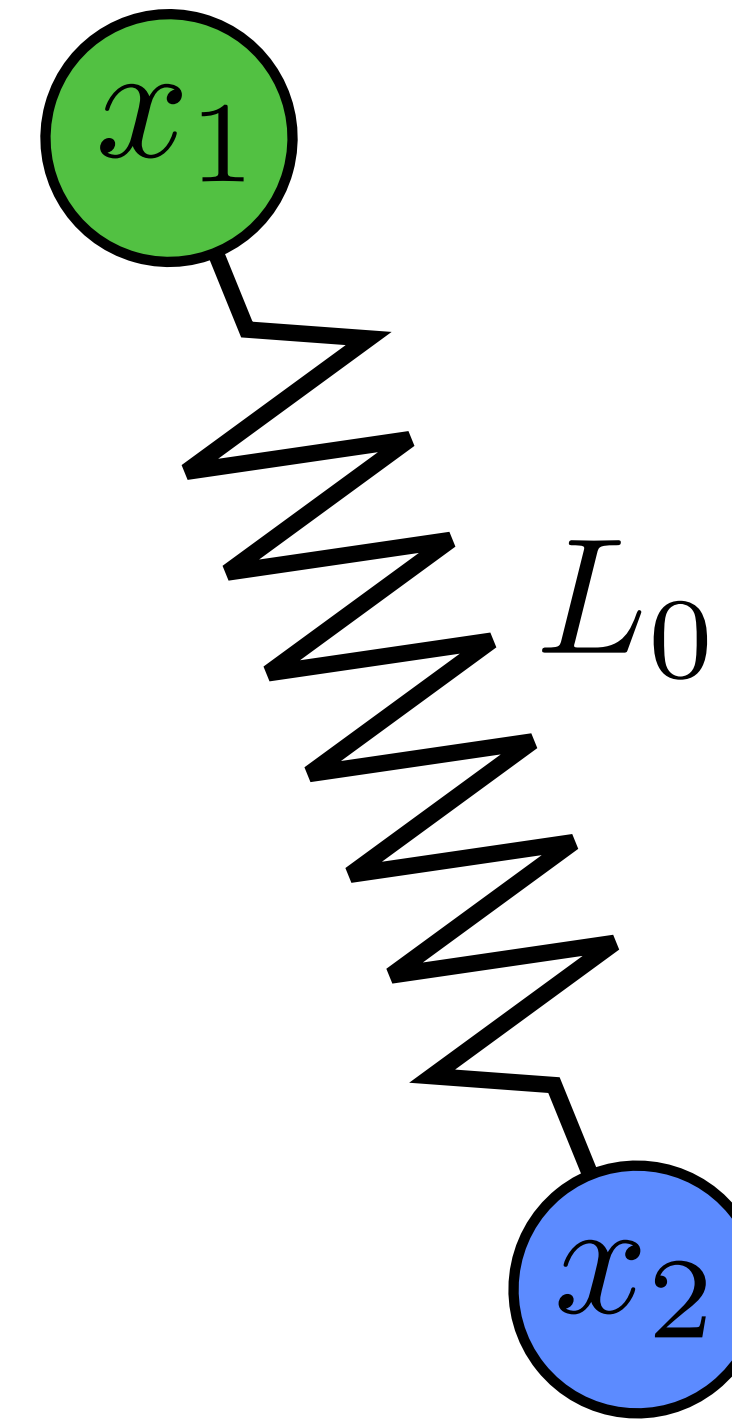
# Example: Mass-Spring System

- Connect particles  $x_1, x_2$  by a spring of length  $L_0$
- Potential energy is given by

$$U = \frac{1}{2}k(L - L_0)^2$$

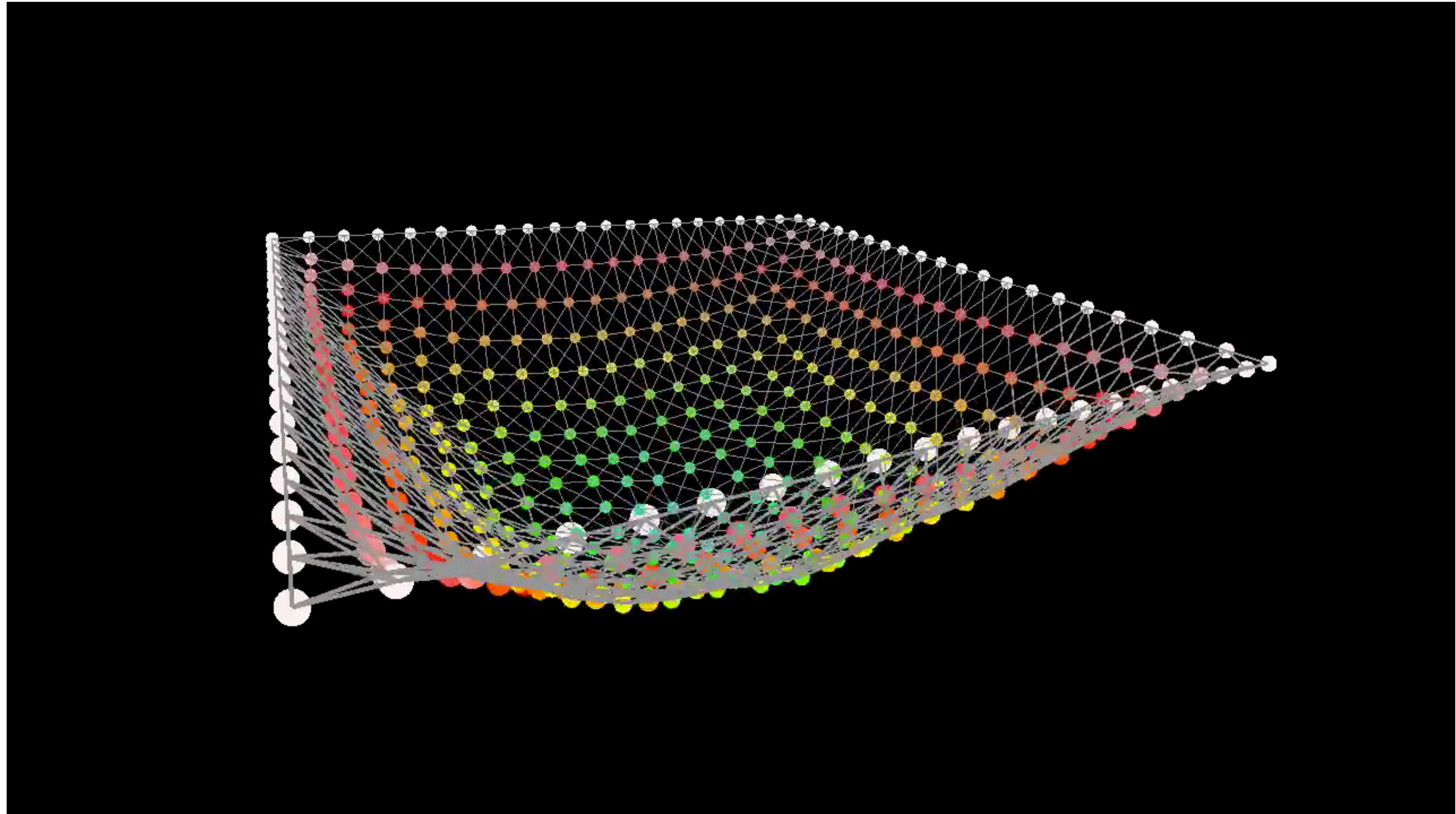
*stiffness*      *current length*      *rest length*

$$= \frac{1}{2}k(|x_1 - x_2|^2 - L_0)^2$$



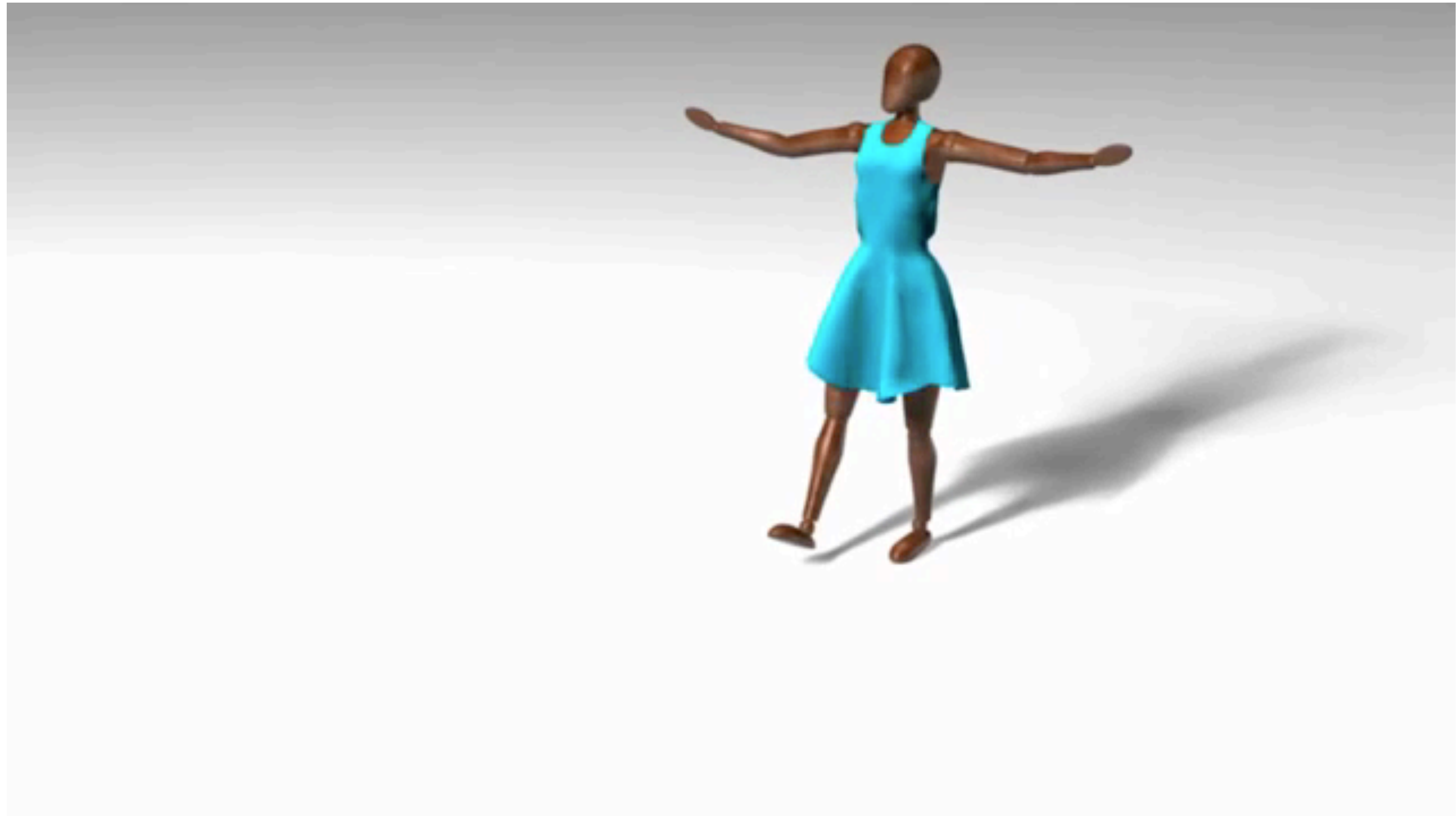
- Connect up many springs to describe interesting phenomena
- Extremely common in graphics/games
  - easy to understand
  - simple equation for each particle
- Often good reasons for using continuum model (PDE)

# Example: Mass Spring System

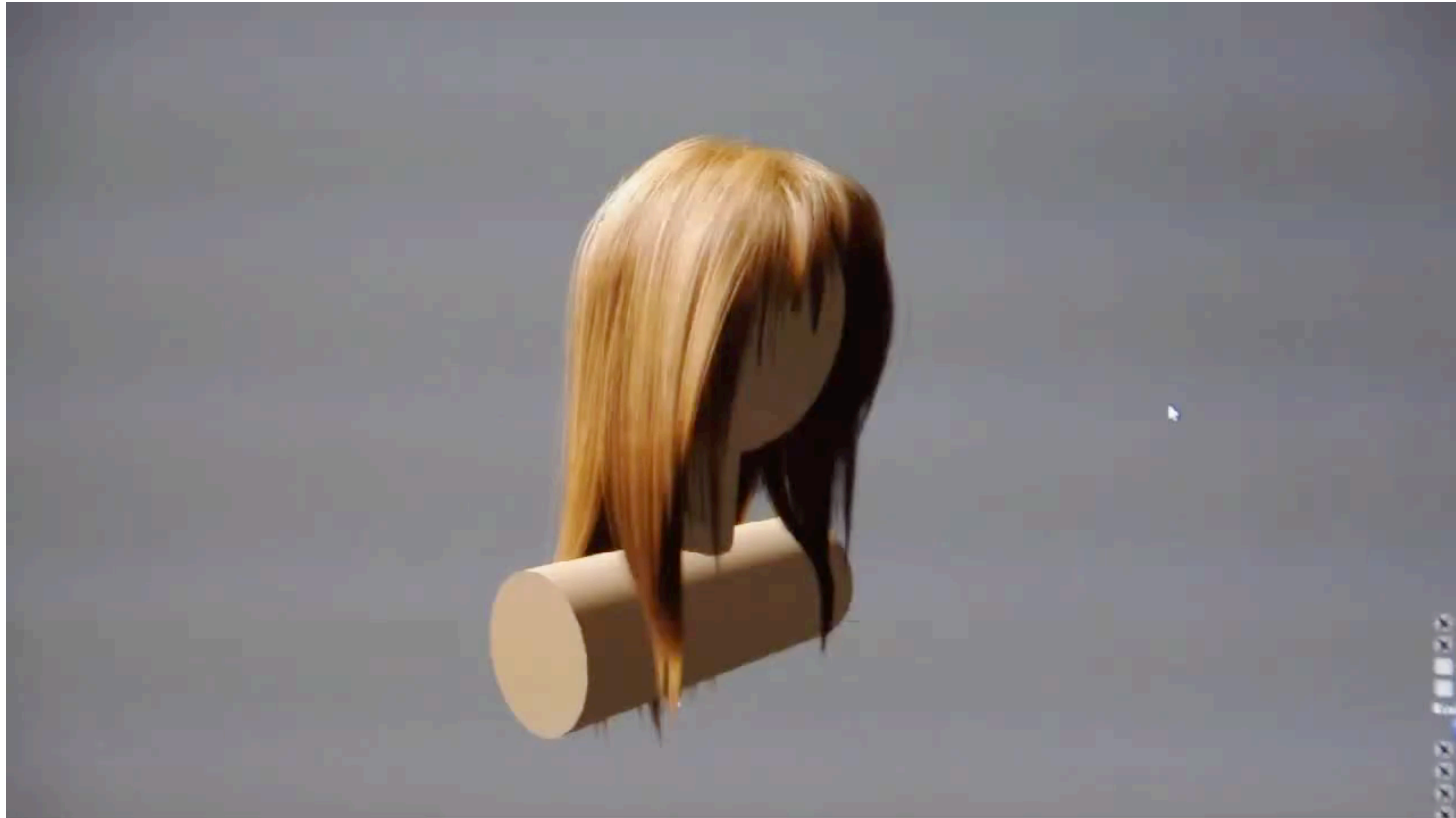




# Example: Mass Spring + Character



# Example: Hair



**Ok, I'm convinced.  
So how do we solve these  
things numerically?**

# Numerical Integration

- Key idea: replace *derivatives* with *differences*
- In ODE, only need to worry about derivative in *time*
- Replace time-continuous function  $q(t)$  with samples  $q_k$  in time

$$\frac{d}{dt} q(t) = f(q(t))$$

⇓

$$\frac{q_{k+1} - q_k}{\tau} = f(q)$$

↙ ↘

**new configuration  
(unknown—want to solve for this!)**

**current configuration  
(known)**

**“time step,” i.e., interval of  
time between  $q_k$  and  $q_{k+1}$**

**Wait... where do we  
evaluate the velocity  
function? At the new  
or old configuration?**

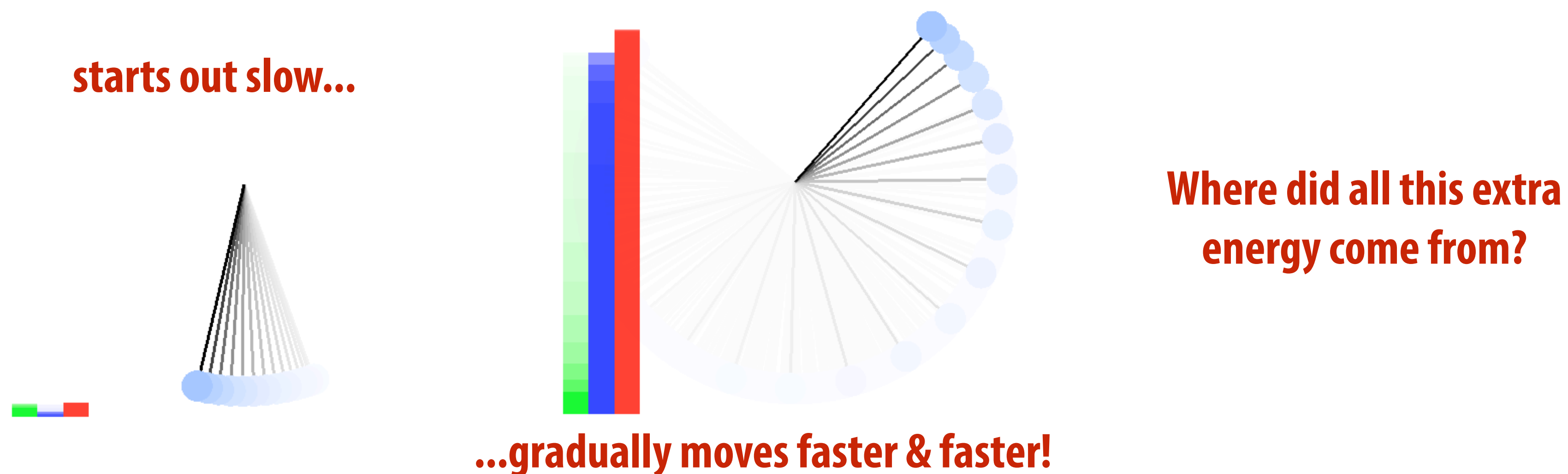
# Forward Euler

- Simplest scheme: evaluate velocity at current configuration
- New configuration can then be written *explicitly* in terms of known data:

$$q_{k+1} = q_k + \tau f(q_k)$$

new configuration      current configuration      velocity at current time

- Very intuitive: walk a tiny bit in the direction of the velocity
- Unfortunately, not very *stable*—consider pendulum:



# Forward Euler - Stability Analysis

- Let's consider behavior of forward Euler for simple linear ODE:

$$\dot{u} = -au, \quad a > 0$$

- Importantly:  $u$  should *decay* (exact solution is  $u(t)=e^{-at}$ )

- Forward Euler approximation is

$$\begin{aligned} u_{k+1} &= u_k - \tau a u_k \\ &= (1 - \tau a) u_k \end{aligned}$$

- Which means after  $n$  steps, we have

$$u_n = (1 - \tau a)^n u_0$$

- Decays only if  $|1-\tau a| < 1$ , or equivalently, if  $\tau < 2/a$
- In practice: need *very small* time steps if  $a$  is large (“stiff system”)

# Backward Euler

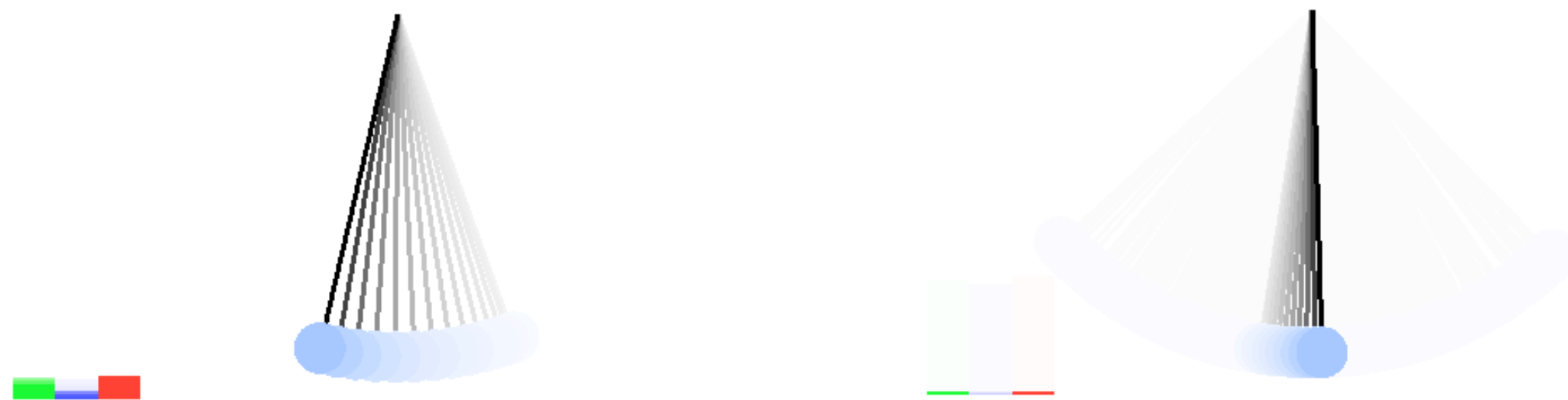
- Let's try something else: evaluate velocity at *next* configuration
- New configuration is then *implicit*, and we must solve for it:

$$q_{k+1} = q_k + \tau f(q_{k+1})$$

new configuration      current configuration      velocity at next time

- Much harder to solve, since in general  $f$  can be very nonlinear!
- Pendulum is now stable... perhaps *too* stable?

starts out slow...



Where did all the energy go?

...and eventually stops moving completely.



# Backward Euler - Stability Analysis

- Again consider a simple linear ODE:

$$\dot{u} = -au, \quad a > 0$$

- Remember:  $u$  should *decay* (exact solution is  $u(t) = e^{-at}$ )

- Backward Euler approximation is

$$(u_{k+1} - u_k) / \tau = -au_{k+1}$$

$$\iff u_{k+1} = \frac{1}{1 + \tau a} u_k$$

- Which means after  $n$  steps, we have

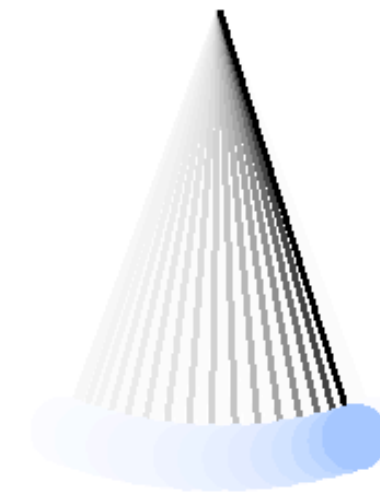
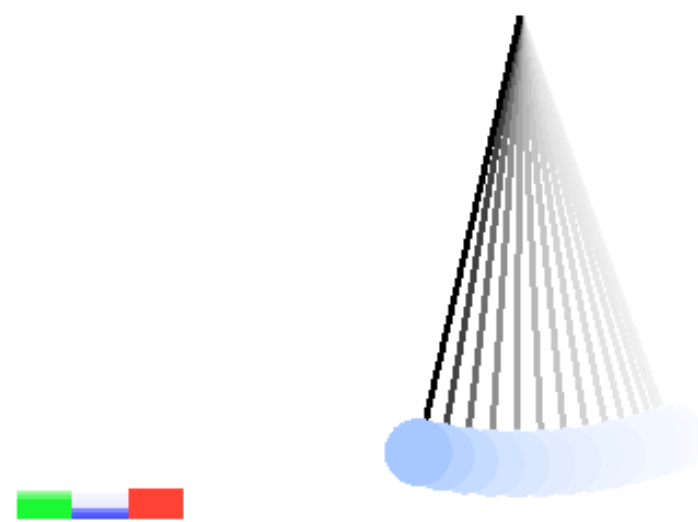
$$u_n = \left( \frac{1}{1 + \tau a} \right)^n u_0$$

- Decays if  $|1 + \tau a| > 1$ , which is always true!
- $\Rightarrow$  Backward Euler is *unconditionally stable* for linear ODEs

# Symplectic Euler

- Backward Euler was stable, but we also saw (empirically) that it exhibits *numerical damping* (damping not found in original eqn.)
- Nice alternative is symplectic Euler
  - update velocity using current configuration
  - update configuration using *new* velocity
- Easy to implement; used often in practice (or leapfrog, Verlet, ...)
- Pendulum now conserves energy *almost exactly*, forever:

starts out slow...

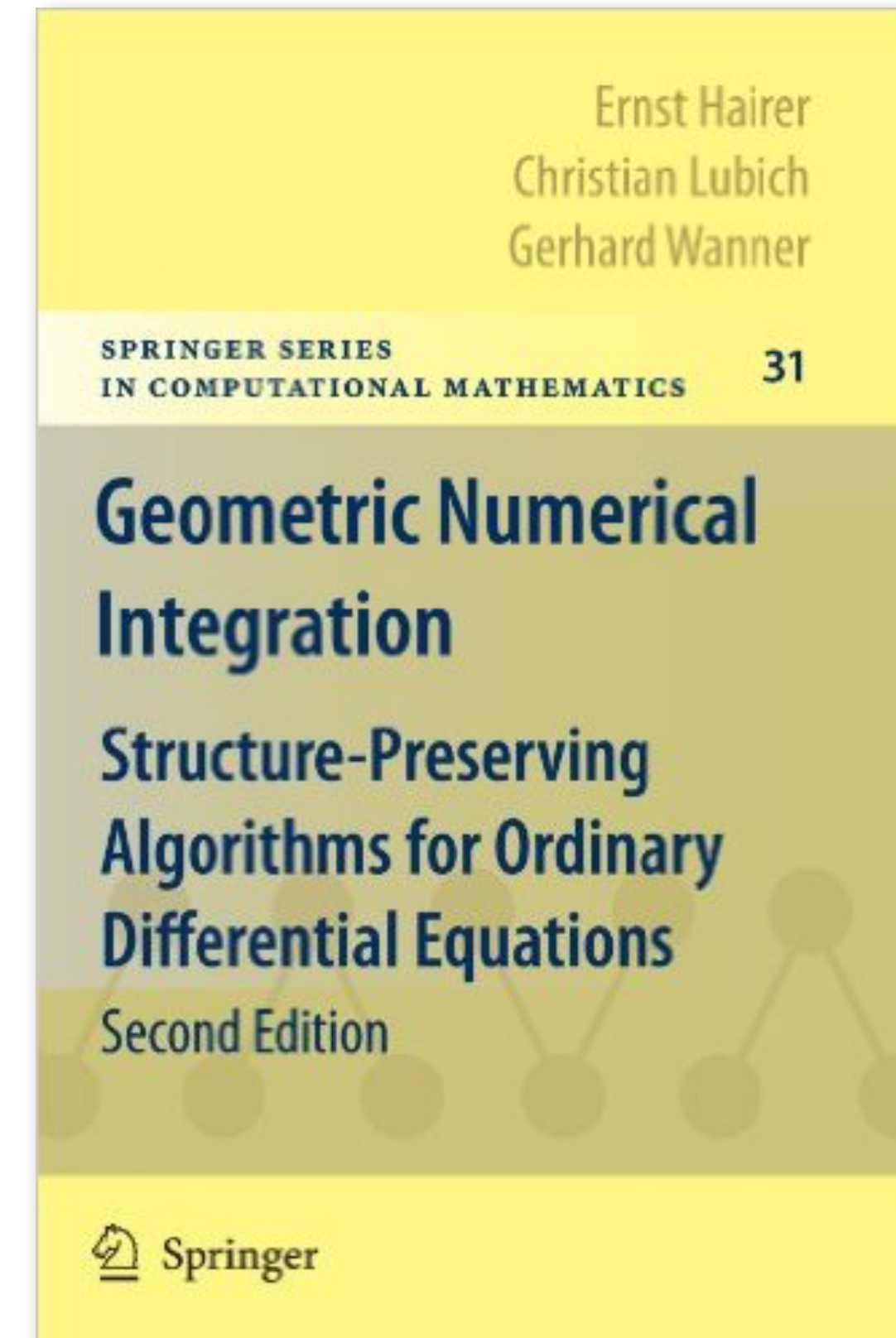


...and keeps on ticking.

(Proof? The analysis is not quite as easy...)

# Numerical Integrators

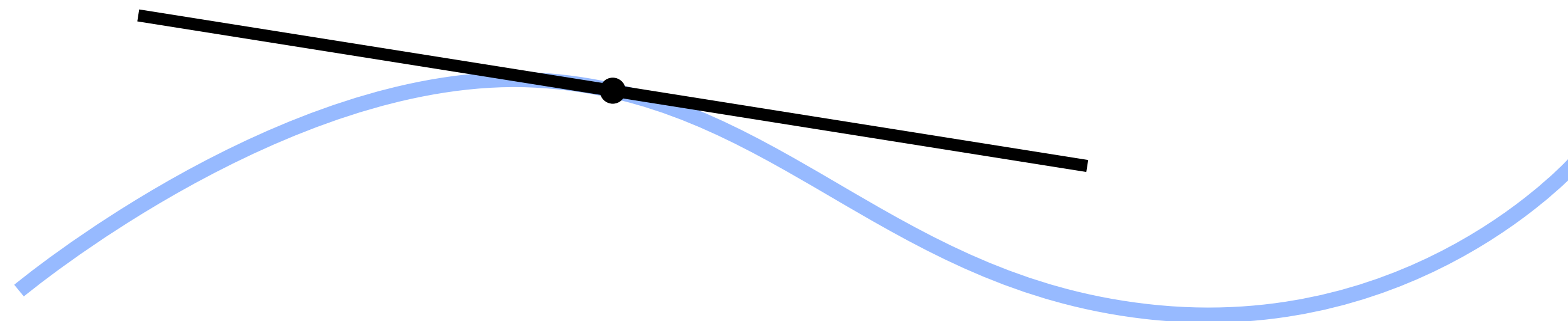
- Barely scratched the surface
- *Many* different integrators
- Why? Because many notions of “good”:
  - stability
  - accuracy
  - consistency/convergence
  - conservation, symmetry, ...
  - computational efficiency (!)
- No one “best” integrator—*pick the right tool for the job!*
- Could do (at least) an entire course on time integration...
- Great book: Hairer, Lubich, Wanner





# Computational Differentiation

- So far, we've been taking derivatives by hand
- Very often in simulation, need to differentiate *extremely complicated* functions (e.g., potential energy, to get forces)
- Several different techniques:
  - keep doing it by hand! (laborious & error prone, but potentially fast)
  - numerical differentiation (simple to code, but usually poor accuracy)
  - automatic differentiation (bigger code investment, better accuracy)
  - symbolic differentiation (can help w/ "by-hand", often messy results)
  - geometric differentiation (sometimes simplifies "by hand" expressions)



# Review: Derivatives

■ Suppose I have a function  $f : \mathbb{R} \rightarrow \mathbb{R}; x \mapsto f(x)$

■ Q: How do I define its first derivative with respect to  $x$ , at  $x_0$ ?

$$f'(x_0) := \lim_{\epsilon \rightarrow 0} \frac{f(x_0 + \epsilon) - f(x_0)}{\epsilon}$$

■ In dynamical simulation, often need to consider functions

$$f : \mathbb{R}^n \rightarrow \mathbb{R}; q \mapsto f(q) \text{ (e.g., potential)}$$

■ *Directional derivative* looks a lot like ordinary derivative:

$$D_X f(q_0) := \lim_{\epsilon \rightarrow 0} \frac{f(q_0 + \epsilon X) - f(q_0)}{\epsilon} \text{ (Q: is } D_X f \text{ vector or scalar?)}$$

■ *Gradient* is vector  $\nabla f$  that yields  $D_X f$  when you take inner product:

$$\langle \nabla f(q_0), X \rangle = D_X f(q_0) \text{ (e.g., gradient of potential is force)}$$

# Numerical Differentiation

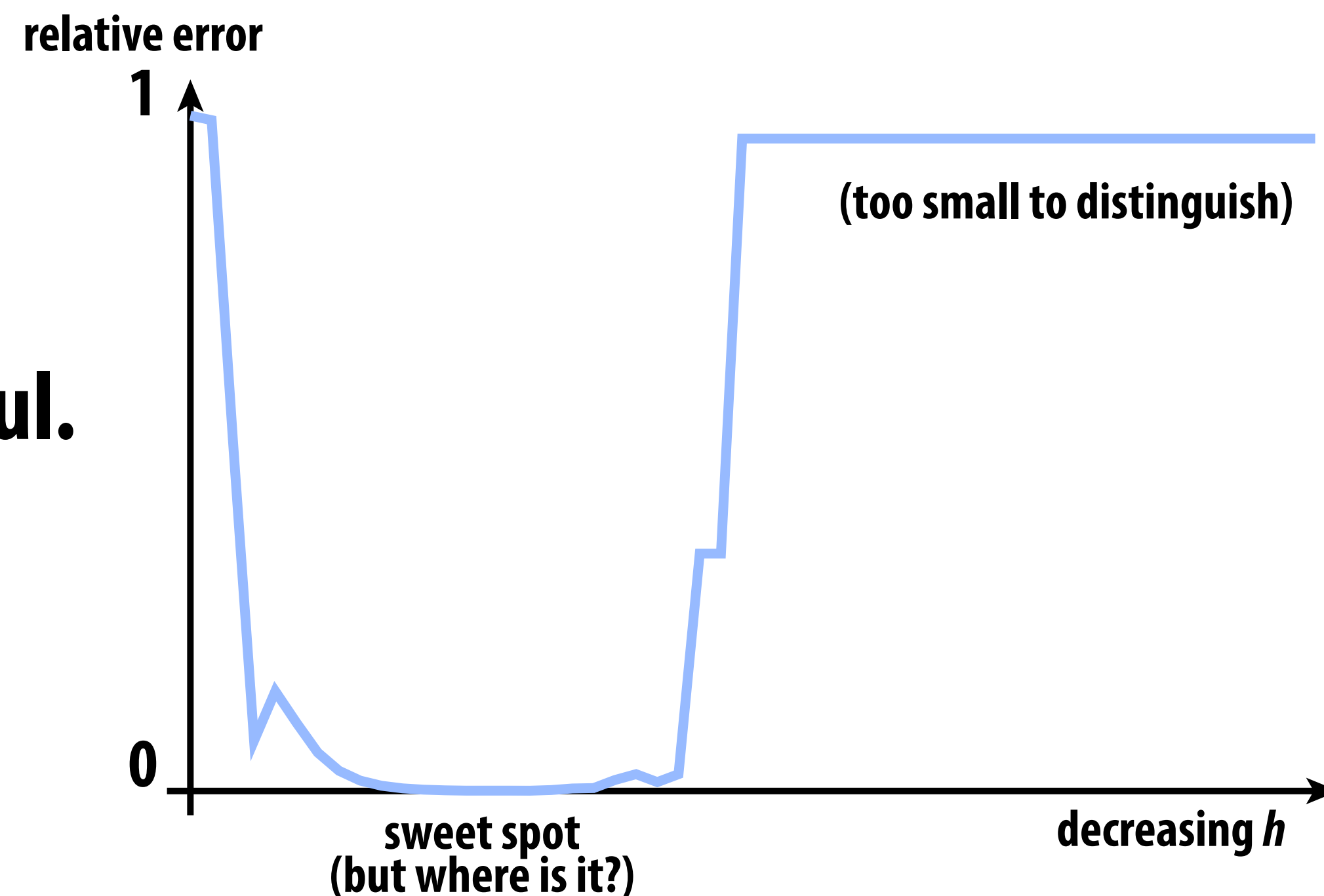
- Taking all those derivatives by hand is a lot of work!  
(Especially if you're just developing/debugging)
- Idea: replace *derivatives* with *differences* (as we did w/ time):

$$f'(x_0) \Rightarrow \frac{f(x_0 + h) - f(x_0)}{h}$$

*h* ← now has *fixed size*

- But how do we pick *h*?
- Smaller is better... right?
- Not always! Must be careful.
- Can also be *expensive!*

e.g., what if *f* were some kind of radiance integral?



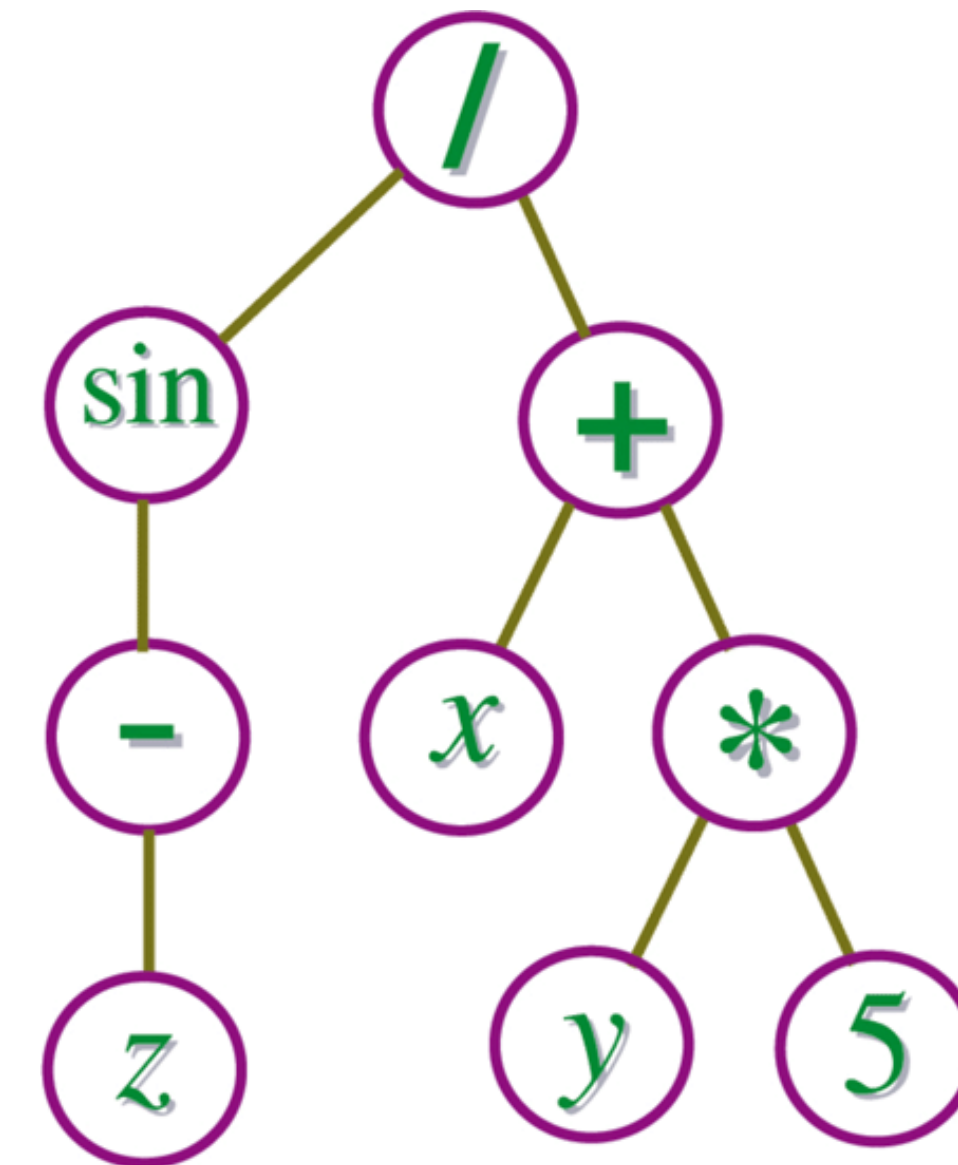


# Automatic Differentiation

- Completely different idea: do arithmetic simultaneously on a function *and* its derivative.
- I.e., rather than work with values  $f$ , work with tuples  $(f, f')$
- Use *chain rule* to determine rules for manipulating tuples
- Example function:  $f(x) = ax^2$
- Suppose we want the value and derivative at  $x=2$
- Start with the tuple  $(x, \frac{\partial}{\partial x} x) |_{x=2} = (2, 1)$
- How do we *multiply* tuples?  $(u, u') * (v, v') = (uv, uv' + vu')$ 
  - for derivatives, we apply the chain rule
  - values just get multiplied
- So, squaring our tuple yields  $(2, 1) * (2, 1) = (4, 4)$ 
  - (did we get it right?)
- And multiplying by  $a$  scales the value *and* derivative:  $(4a, 4a)$
- Pros: good accuracy, *reasonably* fast
  - (must have access to code!)
- Cons: have to redefine all our arithmetic operators!

# Symbolic Differentiation

- Yet another approach (though related to automatic one...)
- Build explicit tree representing expression
- Apply transformations to obtain derivative
- Pros: only needs to happen once!
- Cons: *serious* development investment
- But, can often use existing tools
  - *Mathematica, Maple, etc.*
- Current systems not *great* with vectors, 3D
- Often produce unnecessarily complex formulae...



$$\frac{\sin(-z)}{x + 5y}$$



# Geometric Differentiation

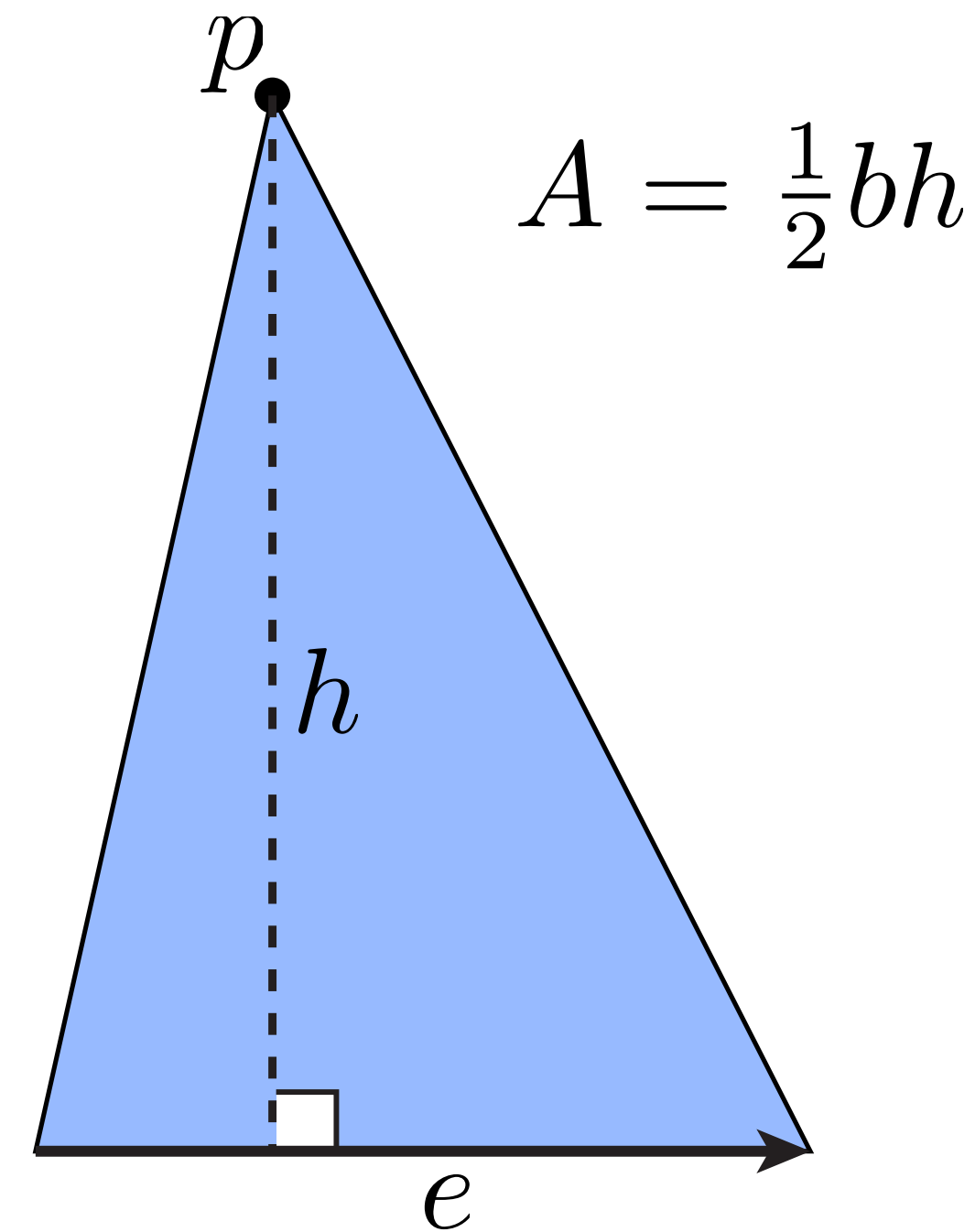
- Sometimes symbolic differentiation misses the “big picture”
- E.g., gradient of triangle area w.r.t. vertex position  $p$

## Mathematica output:

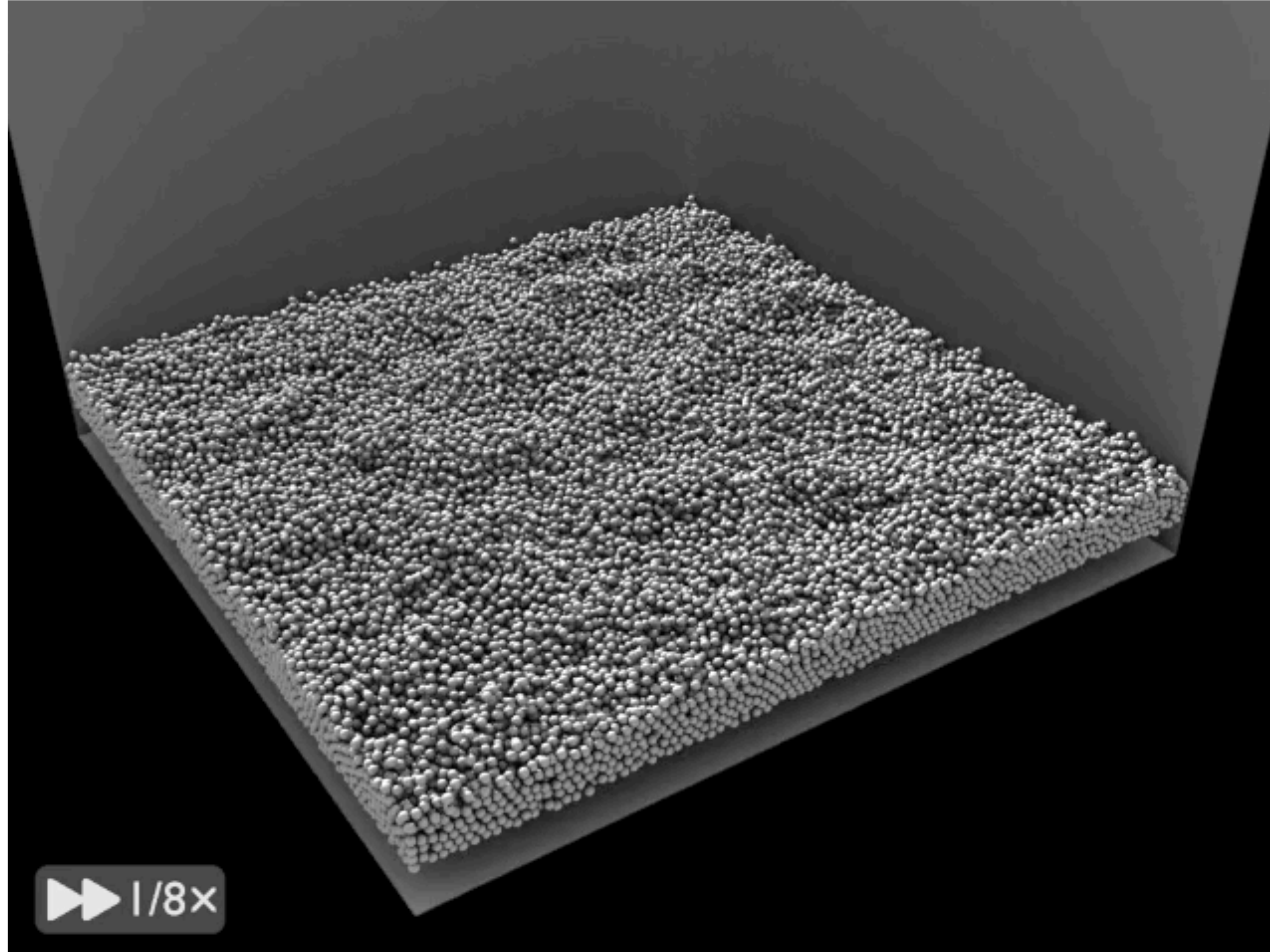
```
(2 (b2 - c2) (-b2 c1 + a2 (-b1 + c1) + a1 (b2 - c2) +
b1 c2) + 2 (b3 - c3) (-b3 c1 + a3 (-b1 + c1) + a1 (b3
- c3) + b1 c3))/(4 Sqrt((a2 b1 - a1 b2 - a2 c1 + b2 c1
+ a1 c2 - p b1 c2)^2 + (a3 b1 - a1 b3 - a3 c1 + b3 c1
+ a1 c3 - b1 c3)^2 + (a3 b2 - a2 b3 - a3 c2 + b3 c2 +
a2 c3 - b2 c3)^2)), (2 (b1 - c1) (a2 (b1 - c1) + b2 c1
- b1 c2 + a1 (-b2 + c2)) + 2 (b3 - c3) (-b3 c2 + a3 (-
b2 + c2) + a2 (b3 - c3) + b2 c3))/(4 Sqrt((a2 b1 - a1
b2 - a2 c1 + b2 c1 + a1 c2 - b1 c2)^2 + (a3 b1 - a1 b3
- a3 c1 + b3 c1 + a1 c3 - b1 c3)^2 + (a3 b2 - a2 b3 -
a3 c2 + b3 c2 + a2 c3 - b2 c3)^2)), (2 (b1 - c1) (a3
(b1 - c1) + b3 c1 - b1 c3 + a1 (-b3 + c3)) + 2 (b2 -
c2) (a3 (b2 - c2) + b3 c2 - b2 c3 + a2 (-b3 + c3)))/(4
Sqrt((a2 b1 - a1 b2 - a2 c1 + b2 c1 + a1 c2 - b1 c2)^2
+ (a3 b1 - a1 b3 - a3 c1 + b3 c1 + a1 c3 - b1 c3)^2 +
(a3 b2 - a2 b3 - a3 c2 + b3 c2 + a2 c3 - b2 c3)^2))
```

## “Geometric” derivative:

$$\nabla_p A = \frac{1}{2} N \times e$$



# Not Covered: Contact Mechanics



Smith et al, *"Reflections on Simultaneous Impact"*



# Coming up next: Optimization

